

Bianchi cosmologies in teleparallel and metric-affine gravity

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 - Relation between gravity and gauge theories?
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- Idea: study geometries with less symmetry:
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 - Bianchi classes II, III, IX admit four-dimensional symmetry groups.
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¹125. anniversary of Bianchi classification!

1. Symmetry in metric-affine and teleparallel geometry
2. The Lorentz algebra
3. Bianchi spacetimes
 - 3.1 General structure
 - 3.2 Bianchi type II
 - 3.3 Bianchi type III
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- Fundamental fields in metric-affine geometry:
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- Metric teleparallelism conventionally formulated using:
 - Tetrad / coframe: $\theta^a = \theta^a{}_{\mu} dx^{\mu}$ with inverse $e_a = e_a{}^{\mu} \partial_{\mu}$.
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Metric-affine and teleparallel geometry

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- Induced metric-affine geometry:

- Metric:

$$g_{\mu\nu} = \eta_{ab} \theta^a{}_\mu \theta^b{}_\nu . \quad (1)$$

- Affine connection:

$$\Gamma^\mu{}_{\nu\rho} = e_a{}^\mu (\partial_\rho \theta^a{}_\nu + \omega^a{}_{b\rho} \theta^b{}_\nu) . \quad (2)$$

Symmetries of metric-affine geometries

- Consider action of Lie algebra \mathfrak{g} spanned by X_A with structure constants $C_{AB}{}^C$:

$$[X_A, X_B] = C_{AB}{}^C X_C. \quad (3)$$

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- Metric:

$$0 = (\mathcal{L}_{X_A} g)_{\mu\nu} = X_A^\rho \partial_\rho g_{\mu\nu} + \partial_\mu X_A^\rho g_{\rho\nu} + \partial_\nu X_A^\rho g_{\mu\rho}. \quad (4)$$

- Connection coefficients:

$$\begin{aligned} 0 &= (\mathcal{L}_{X_A} \Gamma)^\mu{}_{\nu\rho} = X_A^\sigma \partial_\sigma \Gamma^\mu{}_{\nu\rho} - \partial_\sigma X_A^\mu \Gamma^\sigma{}_{\nu\rho} + \partial_\nu X_A^\sigma \Gamma^\mu{}_{\sigma\rho} + \partial_\rho X_A^\sigma \Gamma^\mu{}_{\nu\sigma} + \partial_\nu \partial_\rho X_A^\mu \\ &= \nabla_\rho \nabla_\nu X_A^\mu - X_A^\sigma R^\mu{}_{\nu\rho\sigma} - \nabla_\rho (X_A^\sigma T^\mu{}_{\nu\sigma}). \end{aligned} \quad (5)$$

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- Recall Lie algebra property of vector fields:

$$(\mathcal{L}_{[X_A, X_B]} g)_{\mu\nu} = (\mathcal{L}_{X_A} \mathcal{L}_{X_B} g)_{\mu\nu} - (\mathcal{L}_{X_B} \mathcal{L}_{X_A} g)_{\mu\nu} = 0, \quad (6a)$$

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$$\lambda_{Ac}^a\lambda_{Bb}^c - \lambda_{Bc}^a\lambda_{Ab}^c + X_A^\nu\partial_\nu\lambda_{Bb}^a - X_B^\nu\partial_\nu\lambda_{Ab}^a = C_{AB}{}^C\lambda_{Cb}^a. \quad (9)$$

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$$0 = (\mathcal{L}_{X_A}\omega)^a{}_{b\mu} = \partial_\mu\lambda_{Ab}^a + \omega^a{}_{c\mu}\lambda_{Ab}^c - \omega^c{}_{b\mu}\lambda_{Ac}^a. \quad (11)$$

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$\Rightarrow \lambda_{Ab}^a$ satisfy **property of Lie algebra homomorphism**:

$$\lambda_{Ac}^a\lambda_{Bb}^c - \lambda_{Bc}^a\lambda_{Ab}^c = C_{AB}{}^C\lambda_{Cb}^a. \quad (9)$$

- Metric invariant under X_A^μ if and only if $\lambda_A \in \mathfrak{so}(1, 3)$:

$$\eta_{ac}\lambda_{Ab}^c + \eta_{bc}\lambda_{Aa}^c = 0. \quad (10)$$

- Affine connection is invariant under X_A^μ if and only if $\partial_\mu\lambda_{Ab}^a = 0$:

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- Use Weitzenböck gauge $\omega^a{}_{b\mu} \equiv 0$: **Lie algebra homomorphism : $\mathfrak{g} \rightarrow \mathfrak{so}(1, 3)$** .

- Geometry in Weitzenböck gauge unchanged under global Lorentz transformation Λ^a_b :

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Global Lorentz transformations

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 3. My brain (no idea how much RAM): solved in less than 24 BPU hours.

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- Commutation relations for basis elements:

$$[J_i, J_j] = \epsilon_{ijk} J_k, \quad [K_i, K_j] = -\epsilon_{ijk} J_k, \quad [J_i, K_j] = \epsilon_{ijk} K_k. \quad (16)$$

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- Introduce vector notation for λ_A with $\vec{j}_A, \vec{k}_A \in \mathbb{R}^3$:

$$\lambda_A = \vec{j}_A \cdot \vec{J} + \vec{k}_A \cdot \vec{K}, \quad \vec{J} = (J_1, J_2, J_3), \quad \vec{K} = (K_1, K_2, K_3), \quad (17)$$

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$$[\vec{j} \cdot \vec{J} + \vec{k} \cdot \vec{K}, \vec{j}' \cdot \vec{J} + \vec{k}' \cdot \vec{K}] = (\vec{j} \times \vec{j}' - \vec{k} \times \vec{k}') \cdot \vec{J} + (\vec{j} \times \vec{k}' + \vec{k} \times \vec{j}') \cdot \vec{K}, \quad (19)$$

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⇒ Homomorphism condition for λ :

$$\vec{j}_A \times \vec{j}_B - \vec{k}_A \times \vec{k}_B = C_{AB}^C \vec{j}_C, \quad \vec{j}_A \times \vec{k}_B + \vec{k}_A \times \vec{j}_B = C_{AB}^C \vec{k}_C. \quad (20)$$

Adjoint representation of the Lorentz group

- Adjoint representation of one-parameter groups with unit vector \vec{n} :

$$e^{t\vec{n}\cdot\vec{J}}(\vec{j}\cdot\vec{J})e^{-t\vec{n}\cdot\vec{J}} = [(\vec{n}\cdot\vec{j})\vec{n} + \vec{n}\times\vec{j}\sin t - \vec{n}\times(\vec{n}\times\vec{j})\cos t]\cdot\vec{J}, \quad (21a)$$

$$e^{t\vec{n}\cdot\vec{J}}(\vec{k}\cdot\vec{K})e^{-t\vec{n}\cdot\vec{J}} = [(\vec{n}\cdot\vec{k})\vec{n} + \vec{n}\times\vec{k}\sin t - \vec{n}\times(\vec{n}\times\vec{k})\cos t]\cdot\vec{K}, \quad (21b)$$

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- Adjoint representation of time and space reflections:

$$T\vec{J}T^{-1} = P\vec{J}P^{-1} = \vec{J}, \quad T\vec{K}T^{-1} = P\vec{K}P^{-1} = -\vec{K}. \quad (22)$$

Adjoint representation of the Lorentz group

- Adjoint representation of one-parameter groups with unit vector \vec{n} :

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⇒ Invariant bilinear forms:

$$\kappa(\vec{j}\cdot\vec{J} + \vec{k}\cdot\vec{K}, \vec{j}'\cdot\vec{J} + \vec{k}'\cdot\vec{K}) = \vec{k}\cdot\vec{k}' - \vec{j}\cdot\vec{j}' \quad (23a)$$

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$$S = \vec{j}\cdot\vec{k} = \|\vec{j}\|\|\vec{k}\|\cos\angle(\vec{j}, \vec{k}), \quad D = \|\vec{j}\|^2 - \|\vec{k}\|^2 \quad (24)$$

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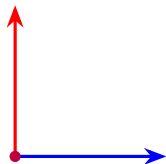
Orbits of the adjoint representation

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 - a. $D = 0$: parametrized by $u \in \mathbb{R}^+$ as $\|\vec{j}\| = \|\vec{k}\| = u > 0$.
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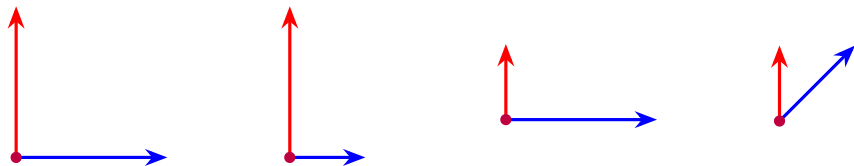


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3. One non-orthogonal orbit for each $|S| > 0$ and $D \in \mathbb{R}$:
 - o Each orbit parametrized by $u = \cos \angle(\vec{j}, \vec{k}) \in [-1, 1] \setminus \{0\}$.
 - o Elements of the orbits satisfy

$$\|\vec{j}\|^2 = \sqrt{\frac{D^2}{4} + \frac{S^2}{u^2}} + \frac{D}{2}, \quad \|\vec{k}\|^2 = \sqrt{\frac{D^2}{4} + \frac{S^2}{u^2}} - \frac{D}{2}. \quad (25)$$

- o Two connected components $\text{sgn } S = \text{sgn } u = \pm 1$ related by reflections.



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2. The Lorentz algebra

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3.2 Bianchi type II

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Free functions in most general geometries

- Number of free functions in most general geometries:

	no symmetry	only homogeneous	1 isotropy	3 isotropies
Metric	$g_{\mu\nu}$	$\mathcal{G}_{1\dots 10}(t)$	$\mathcal{G}_{1\dots 4}(t)$	$\mathcal{G}_{1\dots 2}(t)$
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- Here: consider Bianchi types II, III, IX:
 - General metric-affine geometry: too lengthy for this talk.
 - Metric teleparallel geometry in Weitzenböck gauge: tetrad only.

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- Symmetry algebra:

$$[X_1, X_2] = [X_1, X_3] = [X_1, X_4] = 0, \quad [X_2, X_3] = X_1, \quad [X_2, X_4] = -X_3, \quad [X_3, X_4] = X_2. \quad (26)$$

Symmetry generators

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- Algebra represented by vector fields:

$$X_1 = \partial_y, \quad X_2 = \partial_z, \quad X_3 = z\partial_y - \partial_x, \quad X_4 = z\partial_x + \frac{x^2 - z^2}{2}\partial_y - x\partial_z. \quad (27)$$

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⇒ Homomorphism conditions:

$$\vec{j}_1 \times \vec{j}_2 - \vec{k}_1 \times \vec{k}_2 = 0, \quad \vec{j}_1 \times \vec{k}_2 + \vec{k}_1 \times \vec{j}_2 = 0, \quad (28a)$$

$$\vec{j}_1 \times \vec{j}_3 - \vec{k}_1 \times \vec{k}_3 = 0, \quad \vec{j}_1 \times \vec{k}_3 + \vec{k}_1 \times \vec{j}_3 = 0, \quad (28b)$$

$$\vec{j}_1 \times \vec{j}_4 - \vec{k}_1 \times \vec{k}_4 = 0, \quad \vec{j}_1 \times \vec{k}_4 + \vec{k}_1 \times \vec{j}_4 = 0, \quad (28c)$$

$$\vec{j}_2 \times \vec{j}_3 - \vec{k}_2 \times \vec{k}_3 = \vec{j}_1, \quad \vec{j}_2 \times \vec{k}_3 + \vec{k}_2 \times \vec{j}_3 = \vec{k}_1, \quad (28d)$$

$$\vec{j}_2 \times \vec{j}_4 - \vec{k}_2 \times \vec{k}_4 = -\vec{j}_3, \quad \vec{j}_2 \times \vec{k}_4 + \vec{k}_2 \times \vec{j}_4 = -\vec{k}_3, \quad (28e)$$

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1. Trivial branch: \vec{j}_4, \vec{k}_4 arbitrary and

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2. Non-trivial branch:

○ Conditions to be satisfied:

$$\begin{aligned} \vec{j}_3 \perp \vec{k}_3, \quad \|\vec{j}_3\| = \|\vec{k}_3\|, \quad \vec{j}_4 \perp \vec{k}_4, \quad \|\vec{j}_4\|^2 - \|\vec{k}_4\|^2 = 1, \\ \vec{j}_3 \cdot \vec{k}_4 + \vec{j}_4 \cdot \vec{k}_3 = \vec{j}_3 \cdot \vec{j}_4 - \vec{k}_3 \cdot \vec{k}_4 = 0, \quad \vec{j}_1 = \vec{k}_1 = 0, \\ \vec{j}_3 \times \vec{j}_4 - \vec{k}_3 \times \vec{k}_4 = \vec{j}_2, \quad \vec{j}_3 \times \vec{k}_4 + \vec{k}_3 \times \vec{j}_4 = \vec{k}_2. \end{aligned} \quad (30)$$

⇒ Representative for unique orbit:

$$\vec{j}_4 = \vec{e}_3, \quad \vec{k}_3 = \vec{j}_2 = \vec{e}_1, \quad \vec{j}_3 = -\vec{k}_2 = \vec{e}_2, \quad \vec{j}_1 = \vec{k}_1 = \vec{k}_4 = 0. \quad (31)$$

1. Trivial homomorphism: no non-degenerate tetrad solutions.

Symmetric teleparallel geometry

1. Trivial homomorphism: no non-degenerate tetrad solutions.
2. Non-trivial homomorphism:

- o Tetrad:

$$\theta^0 = [\mathcal{C}_2(1 + x^2 + z^2) + \mathcal{C}_1]dt + (\mathcal{C}_3x + \mathcal{C}_4z)dx + [\mathcal{C}_6(1 + x^2 + z^2) + \mathcal{C}_5]dy + \{[\mathcal{C}_6(1 + x^2 + z^2) + \mathcal{C}_5 - \mathcal{C}_4]x + \mathcal{C}_3z\}dz, \quad (32a)$$

$$\theta^1 = 2\mathcal{C}_2xdt + \mathcal{C}_3dx + 2\mathcal{C}_6xdy + (2\mathcal{C}_6x^2 - \mathcal{C}_4)dz, \quad (32b)$$

$$\theta^2 = 2\mathcal{C}_2zdt + \mathcal{C}_4dx + 2\mathcal{C}_6zdy + (2\mathcal{C}_6xz + \mathcal{C}_3)dz, \quad (32c)$$

$$\theta^3 = [\mathcal{C}_2(1 - x^2 - z^2) - \mathcal{C}_1]dt - (\mathcal{C}_3x + \mathcal{C}_4z)dx + [\mathcal{C}_6(1 - x^2 - z^2) - \mathcal{C}_5]dy + \{[\mathcal{C}_6(1 - x^2 - z^2) - \mathcal{C}_5 + \mathcal{C}_4]x - \mathcal{C}_3z\}dz, \quad (32d)$$

- o Determinant:

$$\det \theta = 2(\mathcal{C}_2\mathcal{C}_5 - \mathcal{C}_1\mathcal{C}_6)(\mathcal{C}_3^2 + \mathcal{C}_4^2). \quad (33)$$

- o Metric:

$$g_{tt} = -4\mathcal{C}_1\mathcal{C}_2, \quad g_{ty} = -2(\mathcal{C}_2\mathcal{C}_5 + \mathcal{C}_1\mathcal{C}_6), \quad g_{tz} = xg_{ty}, \quad g_{xx} = \mathcal{C}_3^2 + \mathcal{C}_4^2, \\ g_{yy} = -4\mathcal{C}_5\mathcal{C}_6, \quad g_{yz} = xg_{yy}, \quad g_{zz} = \mathcal{C}_3^2 + \mathcal{C}_4^2 - 4\mathcal{C}_5\mathcal{C}_6x^2. \quad (34)$$

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$$[X_1, X_2] = [X_2, X_3] = [X_2, X_4] = 0, \quad [X_1, X_3] = -X_1, \quad [X_1, X_4] = X_3, \quad [X_3, X_4] = -X_4. \quad (35)$$

- Algebra represented by vector fields ($0 < n < 1$):

$$X_1 = \partial_y, \quad X_2 = \partial_z, \quad X_3 = \partial_x - y\partial_y, \quad X_4 = y\partial_x + \frac{1}{2} \left(\frac{e^{-2x}}{1-n^2} - y^2 \right) \partial_y - \frac{ne^{-x}}{1-n^2} \partial_z, \quad (36)$$

Symmetry generators

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⇒ Homomorphism conditions:

$$\vec{j}_1 \times \vec{j}_2 - \vec{k}_1 \times \vec{k}_2 = 0, \quad \vec{j}_1 \times \vec{k}_2 + \vec{k}_1 \times \vec{j}_2 = 0, \quad (37a)$$

$$\vec{j}_1 \times \vec{j}_3 - \vec{k}_1 \times \vec{k}_3 = -\vec{j}_1, \quad \vec{j}_1 \times \vec{k}_3 + \vec{k}_1 \times \vec{j}_3 = -\vec{k}_1, \quad (37b)$$

$$\vec{j}_1 \times \vec{j}_4 - \vec{k}_1 \times \vec{k}_4 = \vec{j}_3, \quad \vec{j}_1 \times \vec{k}_4 + \vec{k}_1 \times \vec{j}_4 = \vec{k}_3, \quad (37c)$$

$$\vec{j}_2 \times \vec{j}_3 - \vec{k}_2 \times \vec{k}_3 = 0, \quad \vec{j}_2 \times \vec{k}_3 + \vec{k}_2 \times \vec{j}_3 = 0, \quad (37d)$$

$$\vec{j}_2 \times \vec{j}_4 - \vec{k}_2 \times \vec{k}_4 = 0, \quad \vec{j}_2 \times \vec{k}_4 + \vec{k}_2 \times \vec{j}_4 = 0, \quad (37e)$$

$$\vec{j}_3 \times \vec{j}_4 - \vec{k}_3 \times \vec{k}_4 = -\vec{j}_4, \quad \vec{j}_3 \times \vec{k}_4 + \vec{k}_3 \times \vec{j}_4 = -\vec{k}_4. \quad (37f)$$

1. Trivial branch: \vec{j}_2, \vec{k}_2 arbitrary and

$$\vec{j}_1 = \vec{j}_3 = \vec{j}_4 = \vec{k}_1 = \vec{k}_3 = \vec{k}_4 = 0. \quad (38)$$

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2. Non-trivial branch:

- Conditions to be satisfied:

$$\begin{aligned} \vec{j}_1 \perp \vec{k}_1, \quad \|\vec{j}_1\| &= \|\vec{k}_1\|, \quad \vec{j}_4 \perp \vec{k}_4, \quad \|\vec{j}_4\| = \|\vec{k}_4\|, \\ \vec{j}_1 \cdot \vec{k}_4 + \vec{j}_4 \cdot \vec{k}_1 &= \vec{j}_1 \cdot \vec{j}_4 - \vec{k}_1 \cdot \vec{k}_4 + 1 = 0, \quad \vec{j}_2 = \vec{k}_2 = 0, \\ \vec{j}_1 \times \vec{j}_4 - \vec{k}_1 \times \vec{k}_4 &= \vec{j}_3, \quad \vec{j}_1 \times \vec{k}_4 + \vec{k}_1 \times \vec{j}_4 = \vec{k}_3. \end{aligned} \quad (39)$$

⇒ Representative for unique orbit:

$$\vec{k}_3 = \vec{e}_3, \quad \vec{j}_1 = -\vec{j}_4 = \frac{1}{\sqrt{2}}\vec{e}_1, \quad \vec{k}_1 = \vec{k}_4 = \frac{1}{\sqrt{2}}\vec{e}_2, \quad \vec{j}_2 = \vec{j}_3 = \vec{k}_2 = 0. \quad (40)$$

Symmetric teleparallel geometry

1. Trivial homomorphism: no non-degenerate tetrad solutions.

Symmetric teleparallel geometry

1. Trivial homomorphism: no non-degenerate tetrad solutions.
2. Non-trivial homomorphism (use abbreviation $N = \sqrt{1 - n^2}$):

- o Tetrad:

$$\begin{aligned} \theta^{0,3} &= \frac{C_1}{2\sqrt{2}} \left[\frac{e^{-x}}{N} + N(y^2 \pm 2)e^x \right] dt + \frac{C_4}{2\sqrt{2}n} \left[\frac{e^x}{N} + N(y^2 \pm 2)e^x \right] dz \\ &+ \left\{ \left[\frac{e^{-x}}{N} - (y^2 \pm 2)e^x \right] \frac{C_2}{2\sqrt{2}} - \frac{C_3 y}{\sqrt{2}N} \right\} dx \\ &+ \frac{1}{2\sqrt{2}} \left[N(C_3 + C_4)(y^2 \pm 2)e^{2x} - \frac{C_3 - C_4}{N} - 2ye^x C_2 \right] dy, \end{aligned} \quad (41a)$$

$$\theta^1 = C_5 dt + C_6 e^x dy + \frac{C_6}{n} dz, \quad (41b)$$

$$\theta^2 = -NC_1 ye^x dt + (C_2 ye^x + \frac{C_3}{N}) dx + [C_2 e^x - N(C_3 + C_4) ye^{2x}] dy - \frac{N}{n} C_4 ye^x dz. \quad (41c)$$

- o Determinant:

$$\det \theta = (C_2^2 + C_3^2)(C_4 C_5 - C_1 C_6) \frac{e^x}{nN}. \quad (42)$$

- o Metric:

$$\begin{aligned} g_{tt} &= C_5^2 - C_1^2, & g_{tz} &= \frac{C_5 C_6 - C_1 C_4}{n}, & g_{ty} &= n g_{tz} e^x, & g_{xx} &= \frac{C_2^2 + C_3^2}{N^2}, \\ g_{zz} &= \frac{C_6^2 - C_4^2}{n^2}, & g_{yz} &= n g_{zz} e^x, & g_{yy} &= (C_2^2 + C_3^2 - C_4^2 + C_6^2) e^{2x}. \end{aligned} \quad (43)$$

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4. Conclusion

- Symmetry algebra:

$$[X_1, X_2] = X_3, \quad [X_2, X_3] = X_1, \quad [X_3, X_1] = X_2, \quad [X_1, X_4] = [X_2, X_4] = [X_3, X_4] = 0. \quad (44)$$

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- Algebra represented by vector fields ($n \neq 0$):

$$\begin{aligned} X_1 &= \partial_y, & X_2 &= n \frac{\sin y}{\sin x} \partial_z + \cos y \partial_x - \frac{\sin y}{\tan x} \partial_y, \\ X_3 &= n \frac{\cos y}{\sin x} \partial_z - \sin y \partial_x - \frac{\cos y}{\tan x} \partial_y, & X_4 &= \partial_z, \end{aligned} \quad (45)$$

Symmetry generators

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$$\vec{j}_2 \times \vec{j}_3 - \vec{k}_2 \times \vec{k}_3 = \vec{j}_1, \quad \vec{j}_2 \times \vec{k}_3 + \vec{k}_2 \times \vec{j}_3 = \vec{k}_1, \quad (46a)$$

$$\vec{j}_3 \times \vec{j}_1 - \vec{k}_3 \times \vec{k}_1 = \vec{j}_2, \quad \vec{j}_3 \times \vec{k}_1 + \vec{k}_3 \times \vec{j}_1 = \vec{k}_2, \quad (46b)$$

$$\vec{j}_1 \times \vec{j}_2 - \vec{k}_1 \times \vec{k}_2 = \vec{j}_3, \quad \vec{j}_1 \times \vec{k}_2 + \vec{k}_1 \times \vec{j}_2 = \vec{k}_3, \quad (46c)$$

$$\vec{j}_1 \times \vec{j}_4 - \vec{k}_1 \times \vec{k}_4 = 0, \quad \vec{j}_2 \times \vec{k}_4 + \vec{k}_2 \times \vec{j}_4 = 0, \quad (46d)$$

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2. Non-trivial branch:

○ Conditions to be satisfied:

$$\begin{aligned} \vec{j}_2 \cdot \vec{k}_3 + \vec{j}_3 \cdot \vec{k}_2 = 0, \quad \vec{j}_2 \cdot \vec{j}_3 - \vec{k}_2 \cdot \vec{k}_3 = 0, \quad \|\vec{j}_2\|^2 - \|\vec{k}_2\|^2 = \|\vec{j}_3\|^2 - \|\vec{k}_3\|^2 = 1, \\ \vec{j}_2 \perp \vec{k}_2, \quad \vec{j}_3 \perp \vec{k}_3, \quad \vec{j}_2 \times \vec{j}_3 - \vec{k}_2 \times \vec{k}_3 = \vec{j}_1, \quad \vec{j}_2 \times \vec{k}_3 + \vec{k}_2 \times \vec{j}_3 = \vec{k}_1. \end{aligned} \quad (48)$$

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- o Tetrad:

$$\theta^0 = C_1 dt + nC_2 \cos x dy + C_2 dz, \quad (50a)$$

$$\theta^1 = C_3 \cos x dt - C_4 \sin x dx + (nC_5 \cos^2 x - C_6 \sin^2 x) dy + C_5 \cos x dz, \quad (50b)$$

$$\theta^2 = C_3 \sin x \sin y dt + (C_4 \cos x \sin y - C_6 \cos y) dx + C_5 \sin x \sin y dz \\ + [(nC_5 + C_6) \cos x \sin y + C_4 \cos y] \sin x dy, \quad (50c)$$

$$\theta^3 = C_3 \sin x \cos y dt + (C_4 \cos x \cos y + C_6 \sin y) dx + C_5 \sin x \cos y dz \\ + [(nC_5 + C_6) \cos x \cos y - C_4 \sin y] \sin x dy, \quad (50d)$$

- o Determinant:

$$\det \theta = (C_2 C_3 - C_1 C_5)(C_4^2 + C_6^2) \sin x, \quad (51)$$

- o Metric:

$$g_{tt} = C_3^2 - C_1^2, \quad g_{tz} = C_3 C_5 - C_1 C_2, \quad g_{ty} = n g_{tz} \cos x, \quad g_{xx} = C_4^2 + C_6^2, \\ g_{zz} = C_5^2 - C_2^2, \quad g_{yz} = n g_{zz} \cos x, \quad g_{yy} = g_{xx} \sin^2 x + n^2 g_{zz} \cos^2 x. \quad (52)$$

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- Idea:
 - Study metric-affine and teleparallel geometries with spacetime symmetry algebra \mathfrak{g} .
 - Most general geometry derived from Lie algebra homomorphism $\lambda : \mathfrak{g} \rightarrow \mathfrak{so}(1, 3)$.
 - Equivalent geometries related by adjoint representation of Lorentz transformation.
 - Simple vector product representation to determine most general homomorphisms.
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- Results:
 - Each algebra allows for trivial and non-trivial homomorphisms.
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 - Study coordinate transformations to maximally reduce number of free functions.
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- Reference:

MH, *Spatially homogeneous teleparallel spacetimes with four-dimensional groups of motions*, arXiv:2305.06997 [gr-qc].