

Scalar-torsion theories of gravity

Manuel Hohmann

Laboratory of Theoretical Physics - Institute of Physics - University of Tartu
Center of Excellence “The Dark Side of the Universe”



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- 1 Introduction
- 2 General scalar-torsion gravity
- 3 $L(T, X, Y, \phi)$ theory
- 4 “Scalar-curvature”-like class
- 5 Scalar-torsion gravity without derivative coupling
- 6 Conclusion

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 - Relation between gravity and gauge theories?
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 - Describes gravity as gauge theory of the translation group.
 - Gravitational field strength is torsion.
 - First order action, second order field equations.
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 - Possibly arises from more fundamental theory.
 - Differs from non-minimal coupling to curvature.
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- Scalar field non-minimally coupled to torsion [Geng '11]:
 - Possibly arises from more fundamental theory.
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- Arising questions:
 - Most general class of scalar-torsion gravity theories?
 - Behavior under conformal transformations?

- Fundamental fields:
 - Coframe field $\theta^a = \theta^a{}_\mu dx^\mu$.
 - Flat spin connection $\dot{\omega}^a{}_b = \dot{\omega}^a{}_{b\mu} dx^\mu$.
 - N scalar fields $\phi = (\phi^A; A = 1, \dots, N)$.
 - Arbitrary matter fields χ^I .

- Fundamental fields:

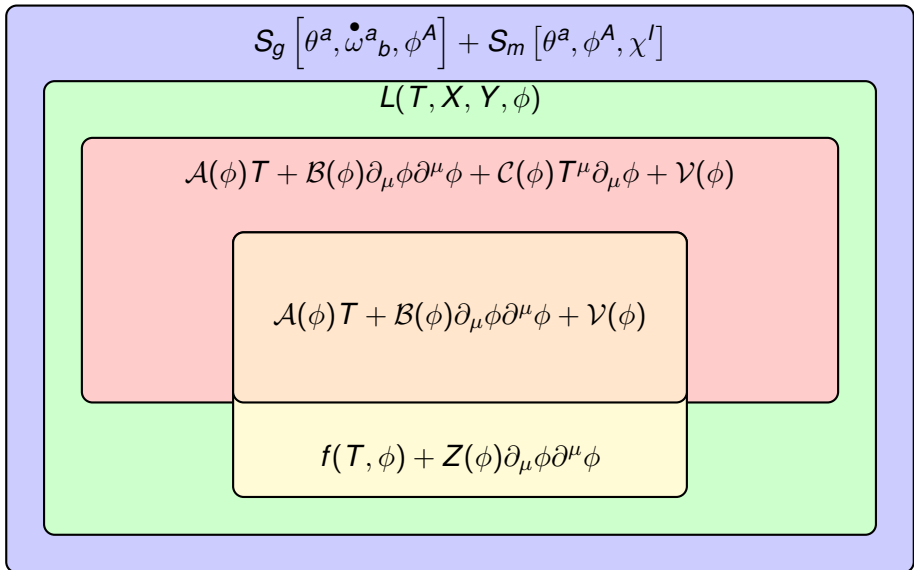
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- Derived quantities:

- Frame field $e_a = e_a{}^\mu \partial_\mu$ with $\iota_{e_a} \theta^b = \delta_a^b$.
- Metric $g_{\mu\nu} = \eta_{ab} \theta^a{}_\mu \theta^b{}_\nu$.
- Volume form $\theta d^4x = \theta^0 \wedge \theta^1 \wedge \theta^2 \wedge \theta^3$.
- Levi-Civita connection

$$\overset{\circ}{\omega}{}_{ab} = -\frac{1}{2}(\iota_{e_b} \iota_{e_c} d\theta_a + \iota_{e_c} \iota_{e_a} d\theta_b - \iota_{e_a} \iota_{e_b} d\theta_c) \theta^c.$$

- Torsion $T^a = d\theta^a + \overset{\circ}{\omega}{}^a{}_b \wedge \theta^b$.



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- Structure of the action [MH '18]:

$$\mathcal{S} \left[\theta^a, \dot{\omega}^a_b, \phi^A, \chi^I \right] = \mathcal{S}_g \left[\theta^a, \dot{\omega}^a_b, \phi^A \right] + \mathcal{S}_m \left[\theta^a, \phi^A, \chi^I \right] .$$

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- Variation of the action:
 - Gravitational part:

$$\begin{aligned} \delta \mathcal{S}_g &= \int_M \left(\Delta_a \wedge \delta \theta^a + \frac{1}{2} \Xi_a{}^b \wedge \delta \dot{\omega}^a{}_b + \Phi_A \wedge \delta \phi^A \right) \\ &= \int_M \left(\Upsilon_a \wedge \delta \theta^a + \Pi_a \wedge \delta T^a + \Phi_A \wedge \delta \phi^A \right) . \end{aligned}$$

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- Matter part:

$$\delta \mathcal{S}_m = \int_M \left(\Sigma_a \wedge \delta \theta^a + \Psi_A \wedge \delta \phi^A + \Omega_I \wedge \delta \chi^I \right) .$$

General scalar-torsion gravity - field equations

- Relation between different terms used to write field equations:

$$\Delta_a = \Upsilon_a - \dot{D}\Pi_a, \quad \Xi^{ab} = -2\Pi^{[a} \wedge \theta^{b]},$$
$$\Pi^a = \frac{1}{4} \iota_{e_c} \iota_{e_b} \Xi^{bc} \wedge \theta^a - \iota_{e_b} \Xi^{ab}, \quad \Upsilon^a = \Delta^a + \dot{D} \left(\frac{1}{4} \iota_{e_c} \iota_{e_b} \Xi^{bc} \wedge \theta^a - \iota_{e_b} \Xi^{ab} \right).$$

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- Scalar field equations: $\Phi_A + \Psi_A = 0$.
- Matter field equations: $\Omega_I = 0$.

- Local Lorentz transformation of the fundamental fields:

$$\delta_\lambda \theta^a = \lambda^a_b \theta^b, \quad \delta_\lambda \dot{\omega}^a_b = \lambda^a_c \dot{\omega}^c_b - \dot{\omega}^a_c \lambda^c_b - d\lambda^a_b = -\dot{D}\lambda^a_b.$$

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$$\delta_\lambda \mathbf{S}_m = \int_M \Sigma_a \wedge (\lambda^a_b \theta^b) = \int_M \Sigma^{[a} \wedge \theta^{b]} \lambda_{ab},$$

$$\begin{aligned} \delta_\lambda \mathbf{S}_g &= \int_M \left[\Upsilon_a \wedge (\lambda^a_b \theta^b) + \Pi_a \wedge (\lambda^a_b T^b) \right] = \int_M \left(\Upsilon^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} \right) \lambda_{ab} \\ &= \int_M \left[\Delta_a \wedge (\lambda^a_b \theta^b) - \frac{1}{2} \Xi_a^b \wedge \dot{D}\lambda^a_b \right] = \int_M \left(\Delta^{[a} \wedge \theta^{b]} - \frac{1}{2} \dot{D}\Xi^{ab} \right) \lambda_{ab}. \end{aligned}$$

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- Consequences:

- Symmetry of the energy-momentum tensor:

$$\Sigma^{[a} \wedge \theta^{b]} = 0.$$

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- Connection equations \equiv antisymmetric part of tetrad equations:

$$\Upsilon^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge \mathcal{T}^{b]} = 0, \quad \Delta^{[a} \wedge \theta^{b]} - \frac{1}{2} \dot{D}\Xi^{ab} = 0.$$

- Variation of the matter action under infinitesimal diffeomorphisms ξ :

$$\delta_{\xi} \mathbf{S}_m = \int_M \left(\Sigma_a \wedge \mathcal{L}_{\xi} \theta^a + \Psi_A \wedge \mathcal{L}_{\xi} \phi^A + \Omega_I \wedge \mathcal{L}_{\xi} \chi^I \right) .$$

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- Energy-momentum conservation:

$$\overset{\circ}{D}\Sigma_a + \Psi_A \wedge \iota_{e_a} d\phi^A = 0.$$

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- Disformal transformation:

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$$\bar{\theta}^a = C(\phi, X)\theta^a + D(\phi, X)\eta^{ab}(\iota_{e_b}\mathbf{d}\phi)\mathbf{d}\phi, \quad X = -\frac{1}{2}\eta^{ab}(\iota_{e_a}\mathbf{d}\phi)(\iota_{e_b}\mathbf{d}\phi)$$

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- Leads to more lengthy relation between original and transformed variables.

- Condition for symmetry under ξ of fundamental fields [MH '15]:

$$\mathcal{L}_\xi \theta^a = -\lambda^a{}_b \theta^b, \quad \mathcal{L}_\xi \dot{\omega}^a{}_b = \dot{D} \lambda^a{}_b, \quad \mathcal{L}_\xi \phi = 0.$$

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- Six generating vector fields ξ_1, \dots, ξ_6 of cosmological symmetry.

- Translation generators:

$$\xi_1 = \chi \sin \vartheta \cos \varphi \partial_r + \frac{\chi}{r} \cos \vartheta \cos \varphi \partial_\vartheta - \frac{\chi \sin \varphi}{r \sin \vartheta},$$

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- Connection field equation solved identically** [MH, L. Järv, M. Krššák, C. Pfeifer '18 to appear].

- Diagonal tetrad in spherical coordinates:

$$\theta^a{}_{\mu} = \text{diag} \left(n(t), \frac{a(t)}{\sqrt{1-kr^2}}, a(t)r, a(t)r \sin \vartheta \right),$$

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- Cosmological spin connections [MH, L. Järv, U. Ualikhanova '18]:

- Spatially flat spacetime $k = 0$:

$$\dot{\omega}^1{}_{2\vartheta} = -\dot{\omega}^2{}_{1\vartheta} = -1, \quad \dot{\omega}^1{}_{3\varphi} = -\dot{\omega}^3{}_{1\varphi} = -\sin \vartheta, \quad \dot{\omega}^2{}_{3\varphi} = -\dot{\omega}^3{}_{2\varphi} = -\cos \vartheta.$$

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- Spatially closed spacetime $k = 1$:

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- Spatially open spacetime $k = -1$:

$$\begin{aligned} \dot{\omega}^0{}_{1r} = \dot{\omega}^1{}_{0r} &= \frac{1}{\sqrt{1+r^2}}, & \dot{\omega}^0{}_{2\vartheta} = \dot{\omega}^2{}_{0\vartheta} &= r, & \dot{\omega}^0{}_{3\varphi} = \dot{\omega}^3{}_{0\varphi} &= r \sin \vartheta, \\ \dot{\omega}^1{}_{2\vartheta} = -\dot{\omega}^2{}_{1\vartheta} &= -\sqrt{1+r^2}, & \dot{\omega}^1{}_{3\varphi} = -\dot{\omega}^3{}_{1\varphi} &= -\sqrt{1+r^2} \sin \vartheta, & \dot{\omega}^2{}_{3\varphi} = -\dot{\omega}^3{}_{2\varphi} &= -\cos \vartheta \end{aligned}$$

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- Gravitational part of the action [MH, C. Pfeifer '18]:

$$S_g[\theta^a, \dot{\omega}^a{}_b, \phi^A] = \int_M L(T, X^{AB}, Y^A, \phi^A) \theta d^4x.$$

- Torsion scalar: $T = \frac{1}{2} T^\rho{}_{\mu\nu} S_\rho{}^{\mu\nu}$.
- Superpotential:

$$S_{\rho\mu\nu} = \frac{1}{2} (T_{\nu\mu\rho} + T_{\rho\mu\nu} - T_{\mu\nu\rho}) - g_{\rho\mu} T^\sigma{}_{\sigma\nu} + g_{\rho\nu} T^\sigma{}_{\sigma\mu}.$$

- Scalar field kinetic term: $X^{AB} = -\frac{1}{2} g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B$.
- Kinetic coupling term: $Y^A = T_\mu{}^{\mu\nu} \phi_{,\nu}^A$.

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- Matter action variation expressed in components:

$$\delta S_m[\theta^a, \phi^A, \chi^I] = \int_M (\Theta_a{}^\mu \delta\theta^a{}_\mu + \vartheta_A \delta\phi^A + \varpi_I \delta\chi^I) \theta d^4x.$$

- Symmetric part of tetrad equations:

$$\begin{aligned} \overset{\circ}{\nabla}_{(\mu} \left(L_{Y^A} \phi_{,\nu)}^A \right) - \overset{\circ}{\nabla}_{\sigma} \left(L_{Y^A} \phi_{,\rho}^A \right) g^{\rho\sigma} g_{\mu\nu} + L_{Y^A} \left(T_{(\mu\nu)}{}^{\rho} \phi_{,\rho}^A + T^{\rho}{}_{\rho(\mu} \phi_{,\nu)}^A \right) \\ - L g_{\mu\nu} - 2 \overset{\circ}{\nabla}_{\rho} \left(L_T \mathcal{S}_{(\mu\nu)}{}^{\rho} \right) + L_T \mathcal{S}_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} - L_{\chi^{AB}} \phi_{,\mu}^A \phi_{,\nu}^B = \Theta_{\mu\nu}. \end{aligned}$$

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- Antisymmetric part of tetrad equations \equiv connection equations:

$$3 \partial_{[\rho} L_T T^{\rho}{}_{\mu\nu]} + \partial_{[\mu} L_{Y^A} \phi_{,\nu]}^A - \frac{3}{2} L_{Y^A} T^{\rho}{}_{[\mu\nu} \phi_{,\rho]}^A = 0.$$

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$$g^{\mu\nu} \overset{\circ}{\nabla}_{\mu} \left(L_{Y^A} T^{\rho}{}_{\rho\nu} - L_{X^{AB}} \phi_{,\nu}^B \right) - L_{\phi^A} = \vartheta_A.$$

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- Matter field equations: $\varpi_I = 0$.

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- Action [MH'18]:

- Gravitational part:

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] = \frac{1}{2\kappa^2} \int_M [-\mathcal{A}(\phi)T + 2\mathcal{B}_{AB}(\phi)X^{AB} + 2\mathcal{C}_A(\phi)Y^A - 2\kappa^2\mathcal{V}(\phi)] \theta d^4x.$$

- Matter part:

$$S_m[\theta^a, \phi^A, \chi^I] = S_m^{\tilde{J}} [e^{\alpha(\phi)}\theta^a, \chi^I].$$

“Scalar-curvature”-like class - action

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- Free functions $\mathcal{A}, \mathcal{B}_{AB}, \mathcal{C}_A, \mathcal{V}, \alpha$ of scalar fields.

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- $\mathcal{C}_A \equiv -\mathcal{A}_{,A} \Leftrightarrow$ theory reduces to scalar-curvature gravity.

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- Free functions $\mathcal{A}, \mathcal{B}_{AB}, \mathcal{C}_A, \mathcal{V}, \alpha$ of scalar fields.
- $\mathcal{C}_A \equiv -\mathcal{A}_{,A} \Leftrightarrow$ theory reduces to scalar-curvature gravity.
- Special subclass of $L(T, X, Y, \phi)$ class of theories.

“Scalar-curvature”-like class - field equation

- Symmetric part of the tetrad equations:

$$\begin{aligned} (\mathcal{A}_{,A} + \mathcal{C}_A) \mathcal{S}_{(\mu\nu)}{}^\rho \phi_{,\rho}^A + \mathcal{A} \left(\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} \overset{\circ}{R} g_{\mu\nu} \right) + \left(\frac{1}{2} \mathcal{B}_{AB} - \mathcal{C}_{(A,B)} \right) \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} \\ - (\mathcal{B}_{AB} - \mathcal{C}_{(A,B)}) \phi_{,\mu}^A \phi_{,\nu}^B + \mathcal{C}_A \left(\overset{\circ}{\nabla}_\mu \overset{\circ}{\nabla}_\nu \phi^A - \overset{\circ}{\square} \phi^A g_{\mu\nu} \right) + \kappa^2 \mathcal{V} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu}, \end{aligned}$$

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- Antisymmetric part of the tetrad field equations:

$$3(\mathcal{A}_{,A} + \mathcal{C}_A) T^\rho{}_{[\mu\nu} \phi_{,\rho]}^A + 2\mathcal{C}_{[A,B]} \phi_{,\mu}^A \phi_{,\nu]}^B = 0.$$

“Scalar-curvature”-like class - field equation

- Symmetric part of the tetrad equations:

$$\begin{aligned}
 (\mathcal{A}_{,A} + C_A) \mathcal{S}_{(\mu\nu)}{}^\rho \phi_{,\rho}^A + \mathcal{A} \left(\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} \overset{\circ}{R} g_{\mu\nu} \right) + \left(\frac{1}{2} \mathcal{B}_{AB} - C_{(A,B)} \right) \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} \\
 - (\mathcal{B}_{AB} - C_{(A,B)}) \phi_{,\mu}^A \phi_{,\nu}^B + C_A \left(\overset{\circ}{\nabla}_\mu \overset{\circ}{\nabla}_\nu \phi^A - \overset{\circ}{\square} \phi^A g_{\mu\nu} \right) + \kappa^2 \mathcal{V} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu},
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- Scalar field equation:

$$\begin{aligned}
 \frac{1}{2} \mathcal{A}_{,A} T - \mathcal{B}_{AB} \overset{\circ}{\square} \phi^B - \left(\mathcal{B}_{AB,C} - \frac{1}{2} \mathcal{B}_{BC,A} \right) g^{\mu\nu} \phi_{,\mu}^B \phi_{,\nu}^C \\
 + C_A \overset{\circ}{\nabla}_\mu T_\nu{}^{\nu\mu} + 2C_{[A,B]} T_\mu{}^{\mu\nu} \phi_{,\nu}^B + \kappa^2 \mathcal{V}_{,A} = \kappa^2 \alpha_{,A} \Theta.
 \end{aligned}$$

“Scalar-curvature”-like class - conformal transf.

- Conformal transformation and scalar field redefinition:

$$\bar{\theta}^a = e^{\gamma(\phi)} \theta^a, \quad \bar{e}_a = e^{-\gamma(\phi)} e_a, \quad \bar{\phi}^A = f^A(\phi).$$

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- Transformation of geometry:

$$\begin{aligned} \bar{T} &= e^{-2\gamma} \left(T + 4\gamma_{,A} Y^A + 12\gamma_{,A}\gamma_{,B} X^{AB} \right), \quad \bar{\phi}^A = f^A, \\ \bar{X}^{AB} &= e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^C} \frac{\partial \bar{\phi}^B}{\partial \phi^D} X^{CD}, \quad \bar{Y}^A = e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^B} \left(Y^B + 6\gamma_{,C} X^{BC} \right), \end{aligned}$$

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- Transformation of parameter functions to preserve action:

$$A = e^{2\gamma} \bar{A},$$

$$B = e^{2\gamma} \left(\bar{B} f'^2 - 6\bar{A} \gamma'^2 + 6\bar{C} f' \gamma' \right),$$

$$C = e^{2\gamma} \left(\bar{C} f' - 2\bar{A} \gamma' \right),$$

$$\mathcal{V} = e^{4\gamma} \bar{\mathcal{V}},$$

$$\alpha = \bar{\alpha} + \gamma.$$

“Scalar-curvature”-like class - invariants

- Quantities invariant under conformal transformations γ :
 - “Scalar” quantities:

$$\mathcal{I}_1 = \frac{e^{2\alpha}}{\mathcal{A}}, \quad \mathcal{I}_2 = \frac{\mathcal{V}}{\mathcal{A}^2}.$$

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$$\mathcal{H}_A = \frac{C_A + \mathcal{A}_{,A}}{2\mathcal{A}}, \quad \mathcal{K}_A = \frac{C_A + 2\alpha_{,A}\mathcal{A}}{2e^{2\alpha}}.$$

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- “Metric” quantities:

$$\mathcal{F}_{AB} = \frac{2\mathcal{A}B_{AB} - 6\mathcal{A}_{,(A}C_{B)} - 3\mathcal{A}_{,A}\mathcal{A}_{,B}}{4\mathcal{A}^2},$$
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$$\mathcal{G}_{AB} = \frac{B_{AB} - 6\alpha_{,(A}C_{B)} - 6\alpha_{,A}\alpha_{,B}\mathcal{A}}{2e^{2\alpha}}.$$

- Covariance under scalar field redefinitions:

$$\bar{\mathcal{I}}_{1,2} = \mathcal{I}_{1,2}, \quad (\bar{\mathcal{H}}, \bar{\mathcal{K}})_A = \frac{\partial\phi^B}{\partial\bar{\phi}^A}(\mathcal{H}, \mathcal{K})_B, \quad (\bar{\mathcal{F}}, \bar{\mathcal{G}})_{AB} = \frac{\partial\phi^C}{\partial\bar{\phi}^A} \frac{\partial\phi^D}{\partial\bar{\phi}^B}(\mathcal{F}, \mathcal{G})_{CD}.$$

- Jordan frame: minimal coupling to matter.

$$\mathcal{A}^{\hat{\mathcal{J}}} = \frac{1}{\mathcal{I}_1}, \quad \mathcal{B}_{AB}^{\hat{\mathcal{J}}} = 2\mathcal{G}_{AB}, \quad \mathcal{C}_A^{\hat{\mathcal{J}}} = 2\mathcal{K}_A, \quad \mathcal{V}^{\hat{\mathcal{J}}} = \frac{\mathcal{I}_2}{\mathcal{I}_1^2}, \quad \alpha^{\hat{\mathcal{J}}} = 0.$$

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- Einstein frame: no coupling to torsion scalar.

$$\mathcal{A}^{\mathfrak{E}} = 1, \quad \mathcal{B}_{AB}^{\mathfrak{E}} = 2\mathcal{F}_{AB}, \quad \mathcal{C}_A^{\mathfrak{E}} = 2\mathcal{H}_A, \quad \mathcal{V}^{\mathfrak{E}} = \mathcal{I}_2, \quad \alpha^{\mathfrak{E}} = \frac{1}{2} \ln \mathcal{I}_1.$$

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- “Debraiding frame” (for $\mathcal{H}_{[A,B]} \equiv 0$): minimal coupling to torsion.

$$\begin{aligned} (\ln \mathcal{A}^{\mathfrak{D}})_{,A} &= 2\mathcal{H}_A, & (\ln \mathcal{B}^{\mathfrak{D}})^A_{B,C} &= [\ln (\mathcal{F} + 3\mathcal{H} \otimes \mathcal{H})]^A_{B,C} + 2\delta_B^A \mathcal{H}_C, \\ \mathcal{C}_A^{\mathfrak{D}} &= 0, & (\ln \mathcal{V}^{\mathfrak{D}})_{,A} &= (\ln \mathcal{I}_2)_A + 4\mathcal{H}_A, & \alpha^{\mathfrak{D}}_{,A} &= \mathcal{I}_1 \mathcal{K}_A. \end{aligned}$$

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- 4 “Scalar-curvature”-like class
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$$S = \frac{1}{2\kappa^2} \int_M \left[f(T, \phi) + Z_{AB}(\phi) g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B \right] \theta d^4x + S_m[\theta^a, \chi^I].$$

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- Field equations:

- Symmetric part of the tetrad field equations:

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- Contains various interesting examples: $f(T)$ equivalent, teleparallel dark energy, ...

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- General scalar-torsion theories:
 - Most general class of theories based on tetrad, flat spin connection, scalar field(s).
 - No direct coupling between matter and spin connection.
 - Local Lorentz invariance: symmetric energy-momentum, dependence of equations.
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- Theory without derivative couplings:
 - Simple, yet interesting class of scalar-torsion theories.
 - Includes $f(T)$, teleparallel dark energy, other studied models.
 - Good test case for application and development of formalisms.

- Cosmological dynamics:
 - Dynamical systems analysis and stable fixed points
 - Dark energy?
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- Hamiltonian formulation:
 - Degrees of freedom in fully dynamical theory.
 - Potential appearance of ghosts?
 - Work towards numerical simulations.

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