

Generalized scalar-torsion theories of gravity

Manuel Hohmann

Laboratory of Theoretical Physics - Institute of Physics - University of Tartu
Center of Excellence "The Dark Side of the Universe"



DPG Spring Conference - Session GR 3 - 20. March 2018

- Open questions in cosmology and gravity:
 - Accelerating phases in the history of the Universe?
 - Relation between gravity and gauge theories?
 - How to quantize gravity?

- Open questions in cosmology and gravity:
 - Accelerating phases in the history of the Universe?
 - Relation between gravity and gauge theories?
 - How to quantize gravity?
- Teleparallel gravity [Møller '61]:
 - Based on tetrad and flat spin connection.
 - Describes gravity as gauge theory of the translation group.
 - Gravitational field strength is torsion.
 - First order action, second order field equations.
 - Spin connection as Lorentz gauge degree of freedom.

- Open questions in cosmology and gravity:
 - Accelerating phases in the history of the Universe?
 - Relation between gravity and gauge theories?
 - How to quantize gravity?
- Teleparallel gravity [Møller '61]:
 - Based on tetrad and flat spin connection.
 - Describes gravity as gauge theory of the translation group.
 - Gravitational field strength is torsion.
 - First order action, second order field equations.
 - Spin connection as Lorentz gauge degree of freedom.
- Scalar field non-minimally coupled to torsion [Geng '11]:
 - Possibly arises from more fundamental theory.
 - Differs from non-minimal coupling to curvature.
 - Possible model for so far unexplained observations.

- Open questions in cosmology and gravity:
 - Accelerating phases in the history of the Universe?
 - Relation between gravity and gauge theories?
 - How to quantize gravity?
- Teleparallel gravity [Møller '61]:
 - Based on tetrad and flat spin connection.
 - Describes gravity as gauge theory of the translation group.
 - Gravitational field strength is torsion.
 - First order action, second order field equations.
 - Spin connection as Lorentz gauge degree of freedom.
- Scalar field non-minimally coupled to torsion [Geng '11]:
 - Possibly arises from more fundamental theory.
 - Differs from non-minimal coupling to curvature.
 - Possible model for so far unexplained observations.
- **Arising questions:**
 - **Most general class of scalar-torsion gravity theories?**
 - **Behavior under conformal transformations?**

- Fundamental fields:

- Coframe field $\theta^a = \theta^a{}_\mu dx^\mu$.
- Flat spin connection $\dot{\omega}^a{}_b = \dot{\omega}^a{}_{b\mu} dx^\mu$.
- N scalar fields $\phi = (\phi^A; A = 1, \dots, N)$.
- Arbitrary matter fields χ^I .

- Fundamental fields:

- Coframe field $\theta^a = \theta^a{}_\mu dx^\mu$.
- Flat spin connection $\overset{\circ}{\omega}{}^a{}_b = \overset{\circ}{\omega}{}^a{}_{b\mu} dx^\mu$.
- N scalar fields $\phi = (\phi^A; A = 1, \dots, N)$.
- Arbitrary matter fields χ^I .

- Derived quantities:

- Frame field $e_a = e_a{}^\mu \partial_\mu$ with $\iota_{e_a} \theta^b = \delta_a^b$.
- Metric $g_{\mu\nu} = \eta_{ab} \theta^a{}_\mu \theta^b{}_\nu$.
- Volume form $\theta d^4x = \theta^0 \wedge \theta^1 \wedge \theta^2 \wedge \theta^3$.
- Levi-Civita connection

$$\overset{\circ}{\omega}{}_{ab} = -\frac{1}{2} (\iota_{e_b} \iota_{e_c} d\theta_a + \iota_{e_c} \iota_{e_a} d\theta_b - \iota_{e_a} \iota_{e_b} d\theta_c) \theta^c.$$

- Torsion $T^a = d\theta^a + \overset{\circ}{\omega}{}^a{}_b \wedge \theta^b$.

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] + S_m [\theta^a, \phi^A, \chi^I]$$

$$L(T, X, Y, \phi)$$

$$A(\phi)T + B(\phi)\partial_\mu\phi\partial^\mu\phi + C(\phi)T^\mu\partial_\mu\phi + V(\phi)$$

$$A(\phi)T + B(\phi)\partial_\mu\phi\partial^\mu\phi + V(\phi)$$

$$f(T, \phi) + Z(\phi)\partial_\mu\phi\partial^\mu\phi$$

- Structure of the action [MH '18]:

$$\mathcal{S} \left[\theta^a, \dot{\omega}^a_b, \phi^A, \chi^I \right] = \mathcal{S}_g \left[\theta^a, \dot{\omega}^a_b, \phi^A \right] + \mathcal{S}_m \left[\theta^a, \phi^A, \chi^I \right] .$$

General scalar-torsion gravity - action

- Structure of the action [MH¹⁸]:

$$S \left[\theta^a, \dot{\omega}^a{}_b, \phi^A, \chi^I \right] = S_g \left[\theta^a, \dot{\omega}^a{}_b, \phi^A \right] + S_m \left[\theta^a, \phi^A, \chi^I \right] .$$

- Variation of the action:
 - Gravitational part:

$$\begin{aligned} \delta S_g &= \int_M \left(\Delta_a \wedge \delta \theta^a + \frac{1}{2} \Xi_a{}^b \wedge \delta \dot{\omega}^a{}_b + \Phi_A \wedge \delta \phi^A \right) \\ &= \int_M \left(\Upsilon_a \wedge \delta \theta^a + \Pi_a \wedge \delta T^a + \Phi_A \wedge \delta \phi^A \right) . \end{aligned}$$

General scalar-torsion gravity - action

- Structure of the action [MH '18]:

$$\mathcal{S} [\theta^a, \dot{\omega}^a{}_b, \phi^A, \chi^I] = \mathcal{S}_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] + \mathcal{S}_m [\theta^a, \phi^A, \chi^I] .$$

- Variation of the action:
 - Gravitational part:

$$\begin{aligned} \delta \mathcal{S}_g &= \int_M \left(\Delta_a \wedge \delta \theta^a + \frac{1}{2} \Xi_a{}^b \wedge \delta \dot{\omega}^a{}_b + \Phi_A \wedge \delta \phi^A \right) \\ &= \int_M \left(\Upsilon_a \wedge \delta \theta^a + \Pi_a \wedge \delta T^a + \Phi_A \wedge \delta \phi^A \right) . \end{aligned}$$

- Matter part:

$$\delta \mathcal{S}_m = \int_M \left(\Sigma_a \wedge \delta \theta^a + \Psi_A \wedge \delta \phi^A + \Omega_I \wedge \delta \chi^I \right) .$$

General scalar-torsion gravity - field equations

- Local Lorentz invariance:

- Gravitational part:

$$\Upsilon^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} = 0 \quad \Leftrightarrow \quad \Delta^{[a} \wedge \theta^{b]} - \frac{1}{2} \dot{D} \Xi^{ab} = 0.$$

- Matter part: $\Sigma^{[a} \wedge \theta^{b]} = 0$.

General scalar-torsion gravity - field equations

- Local Lorentz invariance:

- Gravitational part:

$$\Upsilon^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} = 0 \quad \Leftrightarrow \quad \Delta^{[a} \wedge \theta^{b]} - \frac{1}{2} \overset{\bullet}{D} \Xi^{ab} = 0.$$

- Matter part: $\Sigma^{[a} \wedge \theta^{b]} = 0$.

- Diffeomorphism invariance:

$$\overset{\circ}{D} \Sigma_a + \Psi_A \wedge \iota_{e_a} d\phi^A = 0.$$

General scalar-torsion gravity - field equations

- Local Lorentz invariance:

- Gravitational part:

$$\Upsilon^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} = 0 \quad \Leftrightarrow \quad \Delta^{[a} \wedge \theta^{b]} - \frac{1}{2} \dot{D} \Xi^{ab} = 0.$$

- Matter part: $\Sigma^{[a} \wedge \theta^{b]} = 0$.

- Diffeomorphism invariance:

$$\overset{\circ}{D} \Sigma_a + \Psi_A \wedge \iota_{e_a} d\phi^A = 0.$$

- Field equations:

- Tetrad field equations:

$$\Delta_a + \Sigma_a = 0 \quad \Leftrightarrow \quad \Upsilon_a - \dot{D} \Pi_a + \Sigma_a = 0.$$

- Antisymmetric part \equiv connection field equations:

$$\dot{D} \Xi^{ab} = 0 \quad \Leftrightarrow \quad \dot{D} \Pi^{[a} \wedge \theta^{b]} + \Pi^{[a} \wedge T^{b]} = 0.$$

- Scalar field equations: $\Phi_A + \Psi_A = 0$.

- Matter field equations: $\Omega_I = 0$.

$L(T, X, Y, \phi)$ theory - action

- Gravitational part of the action [MH, C. Pfeifer '18]:

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] = \int_M L(T, X^{AB}, Y^A, \phi^A) \theta d^4x.$$

- Torsion scalar: $T = \frac{1}{2} T^\rho{}_{\mu\nu} S_\rho{}^{\mu\nu}$.
- Superpotential:

$$S_{\rho\mu\nu} = \frac{1}{2} (T_{\nu\mu\rho} + T_{\rho\mu\nu} - T_{\mu\nu\rho}) - g_{\rho\mu} T^\sigma{}_{\sigma\nu} + g_{\rho\nu} T^\sigma{}_{\sigma\mu}.$$

- Scalar field kinetic term: $X^{AB} = -\frac{1}{2} g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B$.
- Kinetic coupling term: $Y^A = T_\mu{}^{\mu\nu} \phi_{,\nu}^A$.

$L(T, X, Y, \phi)$ theory - action

- Gravitational part of the action [MH, C. Pfeifer '18]:

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] = \int_M L(T, X^{AB}, Y^A, \phi^A) \theta d^4x.$$

- Torsion scalar: $T = \frac{1}{2} T^\rho{}_{\mu\nu} S_\rho{}^{\mu\nu}$.
- Superpotential:

$$S_{\rho\mu\nu} = \frac{1}{2} (T_{\nu\mu\rho} + T_{\rho\mu\nu} - T_{\mu\nu\rho}) - g_{\rho\mu} T^\sigma{}_{\sigma\nu} + g_{\rho\nu} T^\sigma{}_{\sigma\mu}.$$

- Scalar field kinetic term: $X^{AB} = -\frac{1}{2} g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B$.
- Kinetic coupling term: $Y^A = T_\mu{}^{\mu\nu} \phi_{,\nu}^A$.
- Matter action variation expressed in components:

$$\delta S_m[\theta^a, \phi^A, \chi^I] = \int_M \left(\Theta_a{}^\mu \delta\theta^a{}_\mu + \vartheta_A \delta\phi^A + \varpi_I \delta\chi^I \right) \theta d^4x.$$

- Symmetric part of tetrad equations:

$$\begin{aligned} & \overset{\circ}{\nabla}_{(\mu} \left(L_{Y^A} \phi_{,\nu)}^A \right) - \overset{\circ}{\nabla}_{\sigma} \left(L_{Y^A} \phi_{,\rho}^A \right) g^{\rho\sigma} g_{\mu\nu} + L_{Y^A} \left(T_{(\mu\nu)}{}^{\rho} \phi_{,\rho}^A + T^{\rho}{}_{\rho(\mu} \phi_{,\nu)}^A \right) \\ & - L g_{\mu\nu} - 2 \overset{\circ}{\nabla}_{\rho} \left(L_T S_{(\mu\nu)}{}^{\rho} \right) + L_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} - L_{X^{AB}} \phi_{,\mu}^A \phi_{,\nu}^B = \Theta_{\mu\nu}. \end{aligned}$$

$L(T, X, Y, \phi)$ theory - field equations

- Symmetric part of tetrad equations:

$$\begin{aligned} \overset{\circ}{\nabla}_{(\mu} (L_{Y^A} \phi_{,\nu}^A) - \overset{\circ}{\nabla}_{\sigma} (L_{Y^A} \phi_{,\rho}^A) g^{\rho\sigma} g_{\mu\nu} + L_{Y^A} (T_{(\mu\nu)}{}^{\rho} \phi_{,\rho}^A + T^{\rho}{}_{\rho(\mu} \phi_{,\nu)}^A) \\ - L g_{\mu\nu} - 2 \overset{\circ}{\nabla}_{\rho} (L_T S_{(\mu\nu)}{}^{\rho}) + L_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} - L_{X^{AB}} \phi_{,\mu}^A \phi_{,\nu}^B = \Theta_{\mu\nu}. \end{aligned}$$

- Antisymmetric part of tetrad equations \equiv connection equations:

$$3 \partial_{[\rho} L_T T^{\rho}{}_{\mu\nu]} + \partial_{[\mu} L_{Y^A} \phi_{,\nu]}^A - \frac{3}{2} L_{Y^A} T^{\rho}{}_{[\mu\nu} \phi_{,\rho]}^A = 0.$$

$L(T, X, Y, \phi)$ theory - field equations

- Symmetric part of tetrad equations:

$$\begin{aligned} \overset{\circ}{\nabla}_{(\mu} \left(L_{Y^A} \phi_{,\nu)}^A \right) - \overset{\circ}{\nabla}_{\sigma} \left(L_{Y^A} \phi_{,\rho}^A \right) g^{\rho\sigma} g_{\mu\nu} + L_{Y^A} \left(T_{(\mu\nu)}{}^{\rho} \phi_{,\rho}^A + T^{\rho}{}_{\rho(\mu} \phi_{,\nu)}^A \right) \\ - L g_{\mu\nu} - 2 \overset{\circ}{\nabla}_{\rho} \left(L_T S_{(\mu\nu)}{}^{\rho} \right) + L_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} - L_{X^{AB}} \phi_{,\mu}^A \phi_{,\nu}^B = \Theta_{\mu\nu}. \end{aligned}$$

- Antisymmetric part of tetrad equations \equiv connection equations:

$$3 \partial_{[\rho} L_T T^{\rho}{}_{\mu\nu]} + \partial_{[\mu} L_{Y^A} \phi_{,\nu]}^A - \frac{3}{2} L_{Y^A} T^{\rho}{}_{[\mu\nu} \phi_{,\rho]}^A = 0.$$

- Scalar field equations:

$$g^{\mu\nu} \overset{\circ}{\nabla}_{\mu} \left(L_{Y^A} T^{\rho}{}_{\rho\nu} - L_{X^{AB}} \phi_{,\nu}^B \right) - L_{\phi^A} = \vartheta_A.$$

$L(T, X, Y, \phi)$ theory - field equations

- Symmetric part of tetrad equations:

$$\begin{aligned} \overset{\circ}{\nabla}_{(\mu} \left(L_{Y^A} \phi_{,\nu)}^A \right) - \overset{\circ}{\nabla}_{\sigma} \left(L_{Y^A} \phi_{,\rho}^A \right) g^{\rho\sigma} g_{\mu\nu} + L_{Y^A} \left(T_{(\mu\nu)}{}^{\rho} \phi_{,\rho}^A + T^{\rho}{}_{\rho(\mu} \phi_{,\nu)}^A \right) \\ - L g_{\mu\nu} - 2 \overset{\circ}{\nabla}_{\rho} \left(L_T S_{(\mu\nu)}{}^{\rho} \right) + L_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} - L_{X^{AB}} \phi_{,\mu}^A \phi_{,\nu}^B = \Theta_{\mu\nu}. \end{aligned}$$

- Antisymmetric part of tetrad equations \equiv connection equations:

$$3 \partial_{[\rho} L_T T^{\rho}{}_{\mu\nu]} + \partial_{[\mu} L_{Y^A} \phi_{,\nu]}^A - \frac{3}{2} L_{Y^A} T^{\rho}{}_{[\mu\nu} \phi_{,\rho]}^A = 0.$$

- Scalar field equations:

$$g^{\mu\nu} \overset{\circ}{\nabla}_{\mu} \left(L_{Y^A} T^{\rho}{}_{\rho\nu} - L_{X^{AB}} \phi_{,\nu}^B \right) - L_{\phi^A} = \vartheta_A.$$

- Matter field equations: $\varpi_I = 0$.

“Scalar-curvature”-like class - action

- Action [MH '18]:

- Gravitational part:

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] = \frac{1}{2\kappa^2} \int_M [-\mathcal{A}(\phi)T + 2\mathcal{B}_{AB}(\phi)X^{AB} + 2\mathcal{C}_A(\phi)Y^A - 2\kappa^2\mathcal{V}(\phi)] \theta d^4x.$$

- Matter part:

$$S_m[\theta^a, \phi, \chi^I] = S_m^\natural [e^{\alpha(\phi)}\theta^a, \chi^I].$$

“Scalar-curvature”-like class - action

- Action [MH '18]:

- Gravitational part:

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] = \frac{1}{2\kappa^2} \int_M [-\mathcal{A}(\phi)T + 2\mathcal{B}_{AB}(\phi)X^{AB} + 2\mathcal{C}_A(\phi)Y^A - 2\kappa^2\mathcal{V}(\phi)] \theta d^4x.$$

- Matter part:

$$S_m[\theta^a, \phi, \chi^I] = S_m^\mathfrak{J} [e^{\alpha(\phi)}\theta^a, \chi^I].$$

- Free functions $\mathcal{A}, \mathcal{B}_{AB}, \mathcal{C}_A, \mathcal{V}, \alpha$ of scalar fields.

“Scalar-curvature”-like class - action

- Action [MH '18]:

- Gravitational part:

$$S_g [\theta^a, \dot{\omega}^a{}_b, \phi^A] = \frac{1}{2\kappa^2} \int_M [-\mathcal{A}(\phi)T + 2\mathcal{B}_{AB}(\phi)X^{AB} + 2\mathcal{C}_A(\phi)Y^A - 2\kappa^2\mathcal{V}(\phi)] \theta d^4x.$$

- Matter part:

$$S_m[\theta^a, \phi, \chi^I] = S_m^\mathcal{J} [e^{\alpha(\phi)}\theta^a, \chi^I].$$

- Free functions $\mathcal{A}, \mathcal{B}_{AB}, \mathcal{C}_A, \mathcal{V}, \alpha$ of scalar fields.
- $\mathcal{C}_a \equiv -\mathcal{A}_{,A} \Leftrightarrow$ theory reduces to scalar-curvature gravity.

“Scalar-curvature”-like class - field equation

- Symmetric part of the tetrad equations:

$$\begin{aligned} & (\mathcal{A}_{,A} + \mathcal{C}_A) \mathcal{S}_{(\mu\nu)}{}^\rho \phi_{,\rho}^A + \mathcal{A} \left(\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} \overset{\circ}{R} g_{\mu\nu} \right) + \left(\frac{1}{2} \mathcal{B}_{AB} - \mathcal{C}_{(A,B)} \right) \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} \\ & - (\mathcal{B}_{AB} - \mathcal{C}_{(A,B)}) \phi_{,\mu}^A \phi_{,\nu}^B + \mathcal{C}_A \left(\overset{\circ}{\nabla}_\mu \overset{\circ}{\nabla}_\nu \phi^A - \overset{\circ}{\square} \phi^A g_{\mu\nu} \right) + \kappa^2 \mathcal{V} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu} , \end{aligned}$$

“Scalar-curvature”-like class - field equation

- Symmetric part of the tetrad equations:

$$\begin{aligned} & (\mathcal{A}_{,A} + \mathcal{C}_A) \mathcal{S}_{(\mu\nu)}{}^\rho \phi_{,\rho}^A + \mathcal{A} \left(\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} \overset{\circ}{R} g_{\mu\nu} \right) + \left(\frac{1}{2} \mathcal{B}_{AB} - \mathcal{C}_{(A,B)} \right) \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} \\ & - (\mathcal{B}_{AB} - \mathcal{C}_{(A,B)}) \phi_{,\mu}^A \phi_{,\nu}^B + \mathcal{C}_A \left(\overset{\circ}{\nabla}_\mu \overset{\circ}{\nabla}_\nu \phi^A - \overset{\circ}{\square} \phi^A g_{\mu\nu} \right) + \kappa^2 \mathcal{V} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu}, \end{aligned}$$

- Antisymmetric part of the tetrad field equations:

$$3(\mathcal{A}_{,A} + \mathcal{C}_A) T^\rho{}_{[\mu\nu]} \phi_{,\rho}^A + 2\mathcal{C}_{[A,B]} \phi_{,\mu}^A \phi_{,\nu}^B = 0.$$

“Scalar-curvature”-like class - field equation

- Symmetric part of the tetrad equations:

$$\begin{aligned}
 & (\mathcal{A}_{,A} + \mathcal{C}_A) \mathcal{S}_{(\mu\nu)}{}^\rho \phi_{,\rho}^A + \mathcal{A} \left(\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} \overset{\circ}{R} g_{\mu\nu} \right) + \left(\frac{1}{2} \mathcal{B}_{AB} - \mathcal{C}_{(A,B)} \right) \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} \\
 & - (\mathcal{B}_{AB} - \mathcal{C}_{(A,B)}) \phi_{,\mu}^A \phi_{,\nu}^B + \mathcal{C}_A \left(\overset{\circ}{\nabla}_\mu \overset{\circ}{\nabla}_\nu \phi^A - \overset{\circ}{\square} \phi^A g_{\mu\nu} \right) + \kappa^2 \mathcal{V} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu},
 \end{aligned}$$

- Antisymmetric part of the tetrad field equations:

$$3(\mathcal{A}_{,A} + \mathcal{C}_A) T^\rho{}_{[\mu\nu]} \phi_{,\rho}^A + 2\mathcal{C}_{[A,B]} \phi_{,\mu}^A \phi_{,\nu}^B = 0.$$

- Scalar field equation:

$$\begin{aligned}
 & \frac{1}{2} \mathcal{A}_{,A} T - \mathcal{B}_{AB} \overset{\circ}{\square} \phi^B - \left(\mathcal{B}_{AB,C} - \frac{1}{2} \mathcal{B}_{BC,A} \right) g^{\mu\nu} \phi_{,\mu}^B \phi_{,\nu}^C \\
 & + \mathcal{C}_A \overset{\circ}{\nabla}_\mu T_\nu{}^{\nu\mu} + 2\mathcal{C}_{[A,B]} T_\mu{}^{\mu\nu} \phi_{,\nu}^B + \kappa^2 \mathcal{V}_{,A} = \kappa^2 \alpha_{,A} \Theta.
 \end{aligned}$$

“Scalar-curvature”-like class - conformal transf.

- Conformal transformation and scalar field redefinition:

$$\bar{\theta}^a = e^{\gamma(\phi)} \theta^a, \quad \bar{e}_a = e^{-\gamma(\phi)} e_a, \quad \bar{\phi}^A = f^A(\phi).$$

“Scalar-curvature”-like class - conformal transf.

- Conformal transformation and scalar field redefinition:

$$\bar{\theta}^a = e^{\gamma(\phi)} \theta^a, \quad \bar{e}_a = e^{-\gamma(\phi)} e_a, \quad \bar{\phi}^A = f^A(\phi).$$

- Transformation of geometry:

$$\begin{aligned} \bar{T} &= e^{-2\gamma} \left(T + 4\gamma_{,A} Y^A + 12\gamma_{,A}\gamma_{,B} X^{AB} \right), \quad \bar{\phi}^A = f^A, \\ \bar{X}^{AB} &= e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^C} \frac{\partial \bar{\phi}^B}{\partial \phi^D} X^{CD}, \quad \bar{Y}^A = e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^B} \left(Y^B + 6\gamma_{,C} X^{BC} \right), \end{aligned}$$

“Scalar-curvature”-like class - conformal transf.

- Conformal transformation and scalar field redefinition:

$$\bar{\theta}^a = e^{\gamma(\phi)} \theta^a, \quad \bar{e}_a = e^{-\gamma(\phi)} e_a, \quad \bar{\phi}^A = f^A(\phi).$$

- Transformation of geometry:

$$\begin{aligned} \bar{T} &= e^{-2\gamma} \left(T + 4\gamma_{,A} Y^A + 12\gamma_{,A}\gamma_{,B} X^{AB} \right), \quad \bar{\phi}^A = f^A, \\ \bar{X}^{AB} &= e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^C} \frac{\partial \bar{\phi}^B}{\partial \phi^D} X^{CD}, \quad \bar{Y}^A = e^{-2\gamma} \frac{\partial \bar{\phi}^A}{\partial \phi^B} \left(Y^B + 6\gamma_{,C} X^{BC} \right), \end{aligned}$$

- Transformation of parameter functions to preserve action:

$$\begin{aligned} \mathcal{A} &= e^{2\gamma} \bar{\mathcal{A}}, \\ \mathcal{B} &= e^{2\gamma} \left(\bar{\mathcal{B}} f'^2 - 6\bar{\mathcal{A}} \gamma'^2 + 6\bar{\mathcal{C}} f' \gamma' \right), \\ \mathcal{C} &= e^{2\gamma} \left(\bar{\mathcal{C}} f' - 2\bar{\mathcal{A}} \gamma' \right), \\ \mathcal{V} &= e^{4\gamma} \bar{\mathcal{V}}, \\ \alpha &= \bar{\alpha} + \gamma. \end{aligned}$$

“Scalar-curvature”-like class - invariants

- Quantities invariant under conformal transformations γ :
 - “Scalar” quantities:

$$\mathcal{I}_1 = \frac{e^{2\alpha}}{\mathcal{A}}, \quad \mathcal{I}_2 = \frac{\mathcal{V}}{\mathcal{A}^2}.$$

- “Covector” quantities:

$$\mathcal{H}_A = \frac{\mathcal{C}_A + \mathcal{A}_{,A}}{2\mathcal{A}}, \quad \mathcal{K}_A = \frac{\mathcal{C}_A + 2\alpha_{,A}\mathcal{A}}{2e^{2\alpha}}.$$

- “Metric” quantities:

$$\mathcal{F}_{AB} = \frac{2\mathcal{A}\mathcal{B}_{AB} - 6\mathcal{A}_{,(A}\mathcal{C}_{B)} - 3\mathcal{A}_{,A}\mathcal{A}_{,B}}{4\mathcal{A}^2},$$

$$\mathcal{G}_{AB} = \frac{\mathcal{B}_{AB} - 6\alpha_{,(A}\mathcal{C}_{B)} - 6\alpha_{,A}\alpha_{,B}\mathcal{A}}{2e^{2\alpha}}.$$

“Scalar-curvature”-like class - invariants

- Quantities invariant under conformal transformations γ :

- “Scalar” quantities:

$$\mathcal{I}_1 = \frac{e^{2\alpha}}{\mathcal{A}}, \quad \mathcal{I}_2 = \frac{\mathcal{V}}{\mathcal{A}^2}.$$

- “Covector” quantities:

$$\mathcal{H}_A = \frac{C_A + \mathcal{A}_{,A}}{2\mathcal{A}}, \quad \mathcal{K}_A = \frac{C_A + 2\alpha_{,A}\mathcal{A}}{2e^{2\alpha}}.$$

- “Metric” quantities:

$$\mathcal{F}_{AB} = \frac{2\mathcal{A}\mathcal{B}_{AB} - 6\mathcal{A}_{,(A}C_{B)} - 3\mathcal{A}_{,A}\mathcal{A}_{,B}}{4\mathcal{A}^2},$$
$$\mathcal{G}_{AB} = \frac{\mathcal{B}_{AB} - 6\alpha_{,(A}C_{B)} - 6\alpha_{,A}\alpha_{,B}\mathcal{A}}{2e^{2\alpha}}.$$

- Covariance under scalar field redefinitions:

$$\bar{\mathcal{I}}_{1,2} = \mathcal{I}_{1,2}, \quad (\bar{\mathcal{H}}, \bar{\mathcal{K}})_A = \frac{\partial\phi^B}{\partial\bar{\phi}^A}(\mathcal{H}, \mathcal{K})_B, \quad (\bar{\mathcal{F}}, \bar{\mathcal{G}})_{AB} = \frac{\partial\phi^C}{\partial\bar{\phi}^A} \frac{\partial\phi^D}{\partial\bar{\phi}^B}(\mathcal{F}, \mathcal{G})_{CD}.$$

“Scalar-curvature”-like class - special frames

- Jordan frame: minimal coupling to matter.

$$\mathcal{A}^{\hat{\mathfrak{J}}} = \frac{1}{\mathcal{I}_1}, \quad \mathcal{B}_{AB}^{\hat{\mathfrak{J}}} = 2\mathcal{G}_{AB}, \quad \mathcal{C}_A^{\hat{\mathfrak{J}}} = 2\mathcal{K}_A, \quad \mathcal{V}^{\hat{\mathfrak{J}}} = \frac{\mathcal{I}_2}{\mathcal{I}_1^2}, \quad \alpha^{\hat{\mathfrak{J}}} = \mathbf{0}.$$

“Scalar-curvature”-like class - special frames

- Jordan frame: minimal coupling to matter.

$$\mathcal{A}^{\mathfrak{J}} = \frac{1}{\mathcal{I}_1}, \quad \mathcal{B}_{AB}^{\mathfrak{J}} = 2\mathcal{G}_{AB}, \quad \mathcal{C}_A^{\mathfrak{J}} = 2\mathcal{K}_A, \quad \mathcal{V}^{\mathfrak{J}} = \frac{\mathcal{I}_2}{\mathcal{I}_1^2}, \quad \alpha^{\mathfrak{J}} = 0.$$

- Einstein frame: no coupling to torsion scalar.

$$\mathcal{A}^{\mathfrak{E}} = 1, \quad \mathcal{B}_{AB}^{\mathfrak{E}} = 2\mathcal{F}_{AB}, \quad \mathcal{C}_A^{\mathfrak{E}} = 2\mathcal{H}_A, \quad \mathcal{V}^{\mathfrak{E}} = \mathcal{I}_2, \quad \alpha^{\mathfrak{E}} = \frac{1}{2} \ln \mathcal{I}_1.$$

“Scalar-curvature”-like class - special frames

- Jordan frame: minimal coupling to matter.

$$\mathcal{A}^{\mathfrak{J}} = \frac{1}{\mathcal{I}_1}, \quad \mathcal{B}_{AB}^{\mathfrak{J}} = 2\mathcal{G}_{AB}, \quad \mathcal{C}_A^{\mathfrak{J}} = 2\mathcal{K}_A, \quad \mathcal{V}^{\mathfrak{J}} = \frac{\mathcal{I}_2}{\mathcal{I}_1^2}, \quad \alpha^{\mathfrak{J}} = 0.$$

- Einstein frame: no coupling to torsion scalar.

$$\mathcal{A}^{\mathfrak{E}} = 1, \quad \mathcal{B}_{AB}^{\mathfrak{E}} = 2\mathcal{F}_{AB}, \quad \mathcal{C}_A^{\mathfrak{E}} = 2\mathcal{H}_A, \quad \mathcal{V}^{\mathfrak{E}} = \mathcal{I}_2, \quad \alpha^{\mathfrak{E}} = \frac{1}{2} \ln \mathcal{I}_1.$$

- “Debraiding frame” (for $\mathcal{H}_{[A,B]} \equiv 0$): minimal coupling to torsion.

$$\begin{aligned} (\ln \mathcal{A}^{\mathfrak{D}})_{,A} &= 2\mathcal{H}_A, & (\ln \mathcal{B}^{\mathfrak{D}})^A_{,B,C} &= [\ln(\mathcal{F} + 3\mathcal{H} \otimes \mathcal{H})]^A_{B,C} + 2\delta_B^A \mathcal{H}_C \\ \mathcal{C}_A^{\mathfrak{D}} &= 0, & (\ln \mathcal{V}^{\mathfrak{D}})_{,A} &= (\ln \mathcal{I}_2)_A + 4\mathcal{H}_A, & \alpha^{\mathfrak{D}}_{,A} &= \mathcal{I}_1 \mathcal{K}_A. \end{aligned}$$

Scalar-torsion gravity without derivative coupling

- Gravitational action [MH, L. Järv, U. Ualikhanova '18]:

$$S = \frac{1}{2\kappa^2} \int_M \left[f(T, \phi) + Z_{AB}(\phi) g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B \right] \theta d^4x + S_m[\theta^a, \chi^I].$$

Scalar-torsion gravity without derivative coupling

- Gravitational action [MH, L. Järv, U. Ualikhanova '18]:

$$S = \frac{1}{2\kappa^2} \int_M \left[f(T, \phi) + Z_{AB}(\phi) g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B \right] \theta d^4x + S_m[\theta^a, \chi^I].$$

- Field equations:

- Symmetric part of the tetrad field equations:

$$\begin{aligned} \frac{1}{2} f g_{\mu\nu} + \overset{\circ}{\nabla}_\rho (f_T S_{(\mu\nu)}{}^\rho) - \frac{1}{2} f_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} \\ - Z_{AB} \phi_{,\mu}^A \phi_{,\nu}^B + \frac{1}{2} Z_{AB} \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu}, \end{aligned}$$

- Antisymmetric part of the tetrad field equations:

$$\partial_{[\rho} f_T T^{\rho}{}_{\mu\nu]} = 0 \quad \Leftrightarrow \quad \partial_\mu f_T \left[\partial_\nu (\theta e_{[a}{}^\mu e_{b]}{}^\nu) + 2\theta e_c{}^{[\mu} e_{[a}{}^{\nu]} \overset{\circ}{\omega}{}^c{}_{b]\nu} \right] = 0.$$

- Scalar field equation:

$$f_{\phi^A} - (2Z_{AB,\phi^C} - Z_{BC,\phi^A}) g^{\mu\nu} \phi_{,\mu}^B \phi_{,\nu}^C - 2Z_{AB} \overset{\circ}{\square} \phi^B = 0.$$

Scalar-torsion gravity without derivative coupling

- Gravitational action [MH, L. Järv, U. Ualikhanova '18]:

$$S = \frac{1}{2\kappa^2} \int_M \left[f(T, \phi) + Z_{AB}(\phi) g^{\mu\nu} \phi_{,\mu}^A \phi_{,\nu}^B \right] \theta d^4x + S_m[\theta^a, \chi^I].$$

- Field equations:

- Symmetric part of the tetrad field equations:

$$\begin{aligned} \frac{1}{2} f g_{\mu\nu} + \overset{\circ}{\nabla}_\rho (f_T S_{(\mu\nu)}{}^\rho) - \frac{1}{2} f_T S_{(\mu}{}^{\rho\sigma} T_{\nu)\rho\sigma} \\ - Z_{AB} \phi_{,\mu}^A \phi_{,\nu}^B + \frac{1}{2} Z_{AB} \phi_{,\rho}^A \phi_{,\sigma}^B g^{\rho\sigma} g_{\mu\nu} = \kappa^2 \Theta_{\mu\nu}, \end{aligned}$$

- Antisymmetric part of the tetrad field equations:

$$\partial_{[\rho} f_T T^{\rho}{}_{\mu\nu]} = 0 \quad \Leftrightarrow \quad \partial_\mu f_T \left[\partial_\nu (\theta e_{[a}{}^\mu e_{b]}{}^\nu) + 2\theta e_c{}^{[\mu} e_{[a}{}^{\nu]} \overset{\circ}{\omega}{}^c{}_{b]\nu} \right] = 0.$$

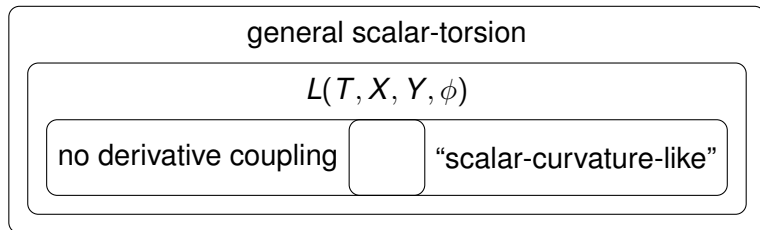
- Scalar field equation:

$$f_{\phi^A} - (2Z_{AB,\phi^C} - Z_{BC,\phi^A}) g^{\mu\nu} \phi_{,\mu}^B \phi_{,\nu}^C - 2Z_{AB} \overset{\circ}{\square} \phi^B = 0.$$

- Solutions: MP 8.3 / Wed, Mar 21, 18:50-19:10, Z6 - SR 1.012.

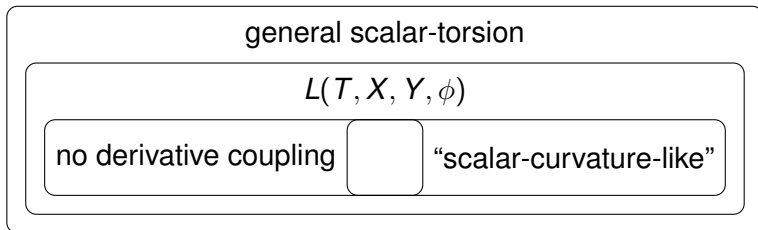
Conclusion

- Summary:
 - Four new classes of scalar-torsion theories.



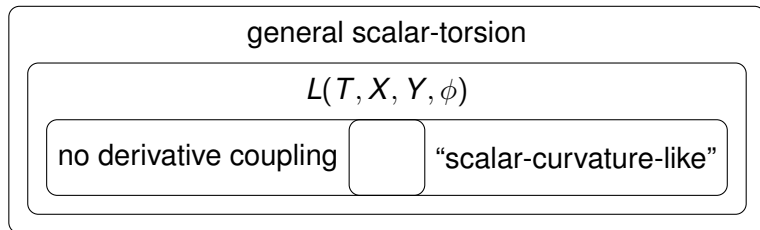
Conclusion

- Summary:
 - Four new classes of scalar-torsion theories.
 - Formulation is covariant under local Lorentz transformations.



Conclusion

- Summary:
 - Four new classes of scalar-torsion theories.
 - Formulation is covariant under local Lorentz transformations.
 - Antisymmetric tetrad field equations \equiv connection field equations.



Conclusion

- Summary:
 - Four new classes of scalar-torsion theories.
 - Formulation is covariant under local Lorentz transformations.
 - Antisymmetric tetrad field equations \equiv connection field equations.
 - **Three classes are closed under conformal transformations.**

general scalar-torsion

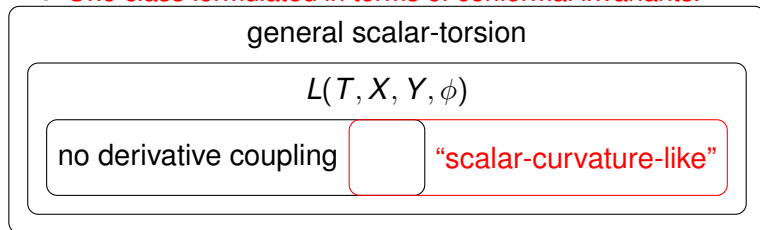
$L(T, X, Y, \phi)$

no derivative coupling

“scalar-curvature-like”

Conclusion

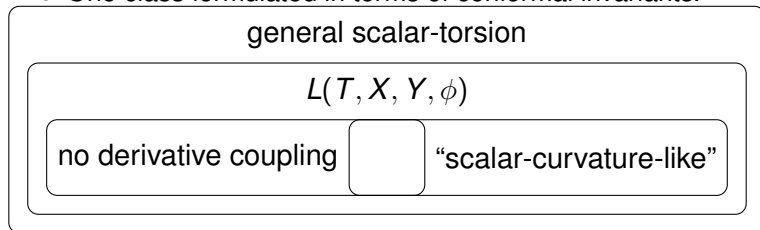
- Summary:
 - Four new classes of scalar-torsion theories.
 - Formulation is covariant under local Lorentz transformations.
 - Antisymmetric tetrad field equations \equiv connection field equations.
 - Three classes are closed under conformal transformations.
 - **One class formulated in terms of conformal invariants.**



Conclusion

- Summary:

- Four new classes of scalar-torsion theories.
- Formulation is covariant under local Lorentz transformations.
- Antisymmetric tetrad field equations \equiv connection field equations.
- Three classes are closed under conformal transformations.
- One class formulated in terms of conformal invariants.



- Outlook - analyze various aspects of these theories:

- Cosmological dynamics / dynamical systems analysis.
- Post-Newtonian limit.
- Gravitational waves - speed and polarisations.
- Hamiltonian formulation - degrees of freedom.

- MH, L. Järv, U. Ualikhanova; Covariant formulation of scalar-torsion gravity; arXiv:1801.05786.
- MH; Scalar-torsion theories of gravity I: general formalism and conformal transformations; arXiv:1801.06528.
- MH, C. Pfeifer; Scalar-torsion theories of gravity II: $L(T, X, Y, \phi)$ theory; arXiv:1801.06536.
- MH; Scalar-torsion theories of gravity III: analogue of scalar-tensor gravity and conformal invariants; arXiv:1801.06531.