

Kinetic theory of fluids in Finsler geometry

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Outline

- 1 Motivation
- 2 Introduction to Finsler spacetimes
- 3 Finslerian description of fluids
- 4 Examples of fluids
- 5 Conclusion

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 - Homogeneity of cosmic microwave background

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- **Idea here: modification of the geometrical structure of spacetime!**
 - **Replace metric spacetime geometry by Finsler geometry.**
 - **Similarly: replacing flat spacetime by curved spacetime led to GR.**

Fluids are everywhere

- Perfect fluid:
 - No shear stress, no friction.
 - Characterized by density ρ and pressure p .
 - Dust, dark matter: $p = 0$.
 - Radiation: $p = \frac{1}{3}\rho$.
 - Dark energy: $p < -\frac{1}{3}\rho$.
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- Charged, multi-component gas:
 - Plasma, interacting gas including recombination / ionization.
 - Used in stellar dynamics, pre-CMB era models. . .

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- Possible explanations of yet unexplained phenomena:
 - Fly-by anomaly
 - Galaxy rotation curves
 - Accelerating expansion of the universe
 - Inflation

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The clock postulate

- Proper time along a curve in Lorentzian spacetime:

$$\tau = \int_{t_1}^{t_2} \sqrt{-g_{ab}(x(t))\dot{x}^a(t)\dot{x}^b(t)} dt .$$

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- Finsler geometry: use a more general length functional:

$$\tau = \int_{t_1}^{t_2} F(x(t), \dot{x}(t)) dt .$$

- Finsler function $F : TM \rightarrow \mathbb{R}^+$.
- Parametrization invariance requires homogeneity:

$$F(x, \lambda y) = \lambda F(x, y) \quad \forall \lambda > 0 .$$

Finsler spacetime geometry

- Finsler geometries suitable for spacetimes exist. [C. Pfeifer, M. Wohlfarth '11]

⇒ Finsler metric with Lorentz signature:

$$g_{ab}^F(x, y) = \frac{1}{2} \frac{\partial}{\partial y^a} \frac{\partial}{\partial y^b} F^2(x, y).$$

⇒ Notion of timelike, lightlike, spacelike tangent vectors.

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- Unit vectors $y \in T_x M$ defined by

$$F^2(x, y) = g_{ab}^F(x, y) y^a y^b = 1.$$

⇒ Set $\Omega_x \subset T_x M$ of unit timelike vectors at $x \in M$.

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- Ω_x contains a closed connected component $S_x \subseteq \Omega_x$.

↪ **Causality: S_x corresponds to physical observers.**

Geometry on the tangent bundle

- Cartan non-linear connection:

$$N^a_b = \frac{1}{4} \bar{\partial}_b \left[g^{F ac} (y^d \partial_d \bar{\partial}_c F^2 - \partial_c F^2) \right]$$

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⇒ Split of the tangent and cotangent bundles:

- Tangent bundle: $TTM = HTM \oplus VTM$

$$\delta_a = \partial_a - N^b_a \bar{\partial}_b, \quad \bar{\partial}_a$$

- Cotangent bundle: $T^*TM = H^*TM \oplus V^*TM$

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$$G = -g_{ab}^F dx^a \otimes dx^b - \frac{g_{ab}^F}{F^2} \delta y^a \otimes \delta y^b$$

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- Geodesic spray:

$$\mathbf{S} = y^a \delta_a$$

- Recall from the definition of Finsler spacetimes:
 - Set $\Omega_x \subset T_x M$ of unit timelike vectors at $x \in M$.
 - Physical observers correspond to $S_x \subseteq \Omega_x$.
- Definition of observer space:

$$O = \bigcup_{x \in M} S_x \subset TM.$$

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- Sasaki metric \tilde{G} on O given by pullback of G to O .
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- Geodesic hypersurface measure $\omega = \iota_{\mathbf{r}} \Sigma$.
- Note that $\mathcal{L}_{\mathbf{r}} \Sigma = 0$ and $d\omega = 0$.

From metric to Finsler geometry

Tangent bundle geometry:

- ○ Finsler function:

$$F(x, y) = \sqrt{|g_{ab}(x)y^a y^b|}$$

- Finsler metric:

$$g_{ab}^F(x, y) = \begin{cases} -g_{ab}(x) & y \text{ timelike} \\ g_{ab}(x) & y \text{ spacelike} \end{cases}$$

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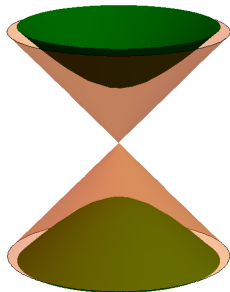
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- Observer space:

- Space Ω_x of unit timelike vectors at $x \in M$.
- Space S_x of future unit timelike vectors at $x \in M$.
- Observer space O : union of shells S_x .

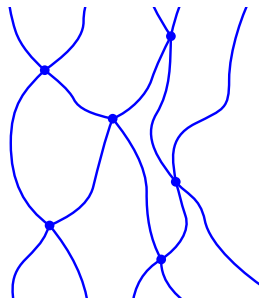


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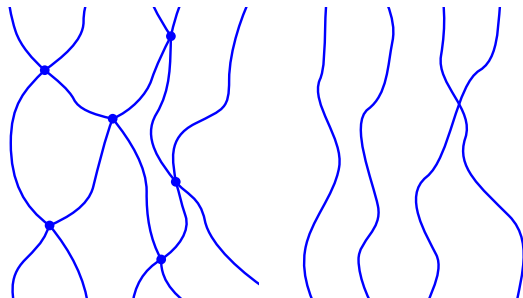
Definition of fluids

- Single-component fluid:
 - Constituted by classical, relativistic particles.
 - Particles have equal properties (mass, charge, ...).
 - Particles follow piecewise geodesic curves.
 - Endpoints of geodesics are interactions with other particles.



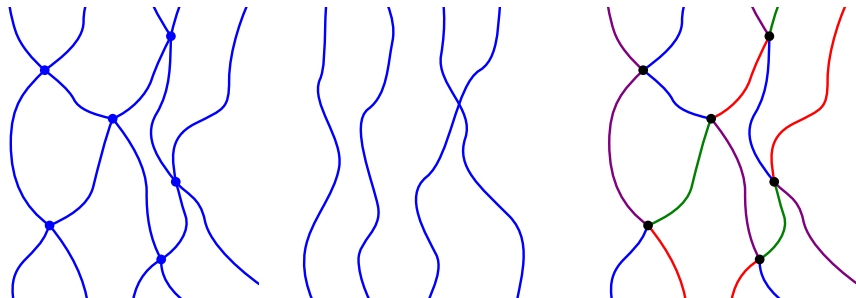
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- **Multi-component fluid: multiple types of particles.**



Geodesics on observer space

- Dynamics of fluids depends on geodesic equation.
- Geodesic equation for curve $x(\tau)$ on spacetime M :

$$\ddot{x}^a + N^a_b(x, \dot{x})\dot{x}^b = 0.$$

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$$\dot{x}^a = y^a, \quad \dot{y}^a = -N^a_b(x, y)y^b.$$

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- Tangent vectors are future unit timelike: $(x, y) \in O$.
- ⇒ Particle trajectories are piecewise integral curves of \mathbf{r} on O .

One-particle distribution function

- Recall: $\omega = \iota_{\mathbf{r}}\Sigma \in \Omega^6(\mathcal{O})$ unique 6-form such that:
 - ω non-degenerate on every hypersurface not tangent to \mathbf{r} .
 - $d\omega = 0$.

One-particle distribution function

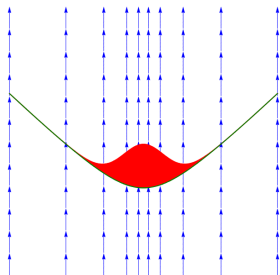
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For every hypersurface $\sigma \subset \mathcal{O}$,

$$N[\sigma] = \int_{\sigma} \phi \omega$$

of **particle trajectories** through σ .

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- Counting of particle trajectories respects hypersurface orientation.



One-particle distribution function

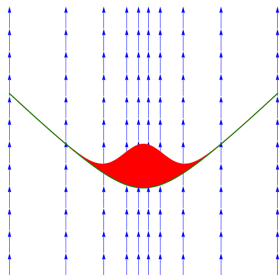
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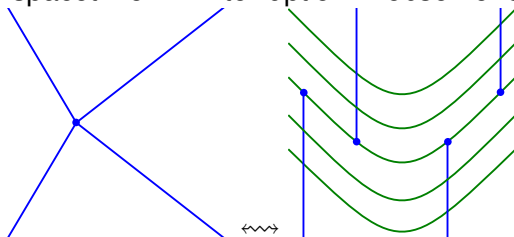
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- For multi-component fluids: ϕ_i for each component i .



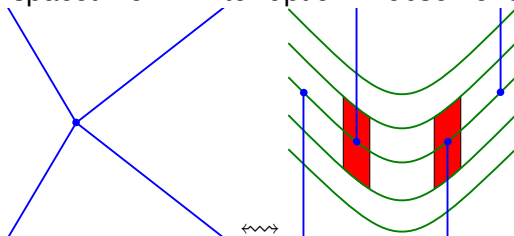
Collisions & the Liouville equation

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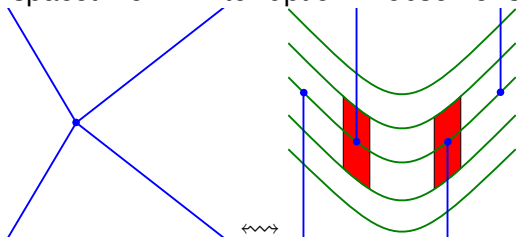
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of outbound trajectories - # of inbound trajectories.

\Rightarrow Collision density measured by $\mathcal{L}_r \phi$.

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- Collisionless fluid: trajectories have no endpoints, $\mathcal{L}_{\mathbf{r}} \phi = 0$.**

\Rightarrow Simple, first order equation of motion for collisionless fluid.

\Rightarrow ϕ is constant along integral curves of \mathbf{r} .

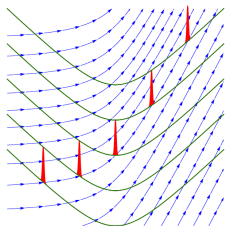
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Some (very) pictorial examples

Geodesic dust fluid:

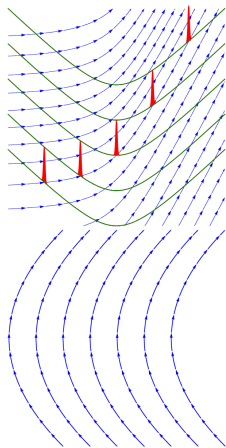
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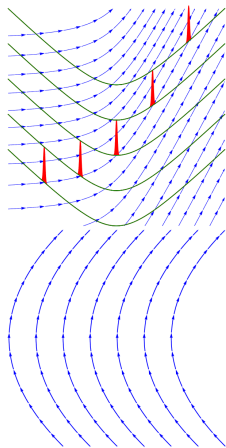


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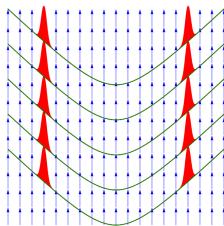
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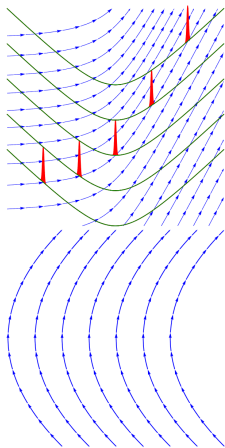
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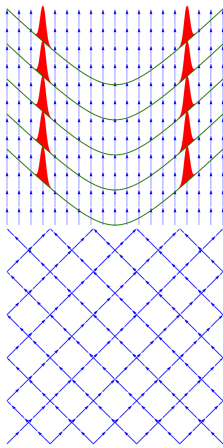
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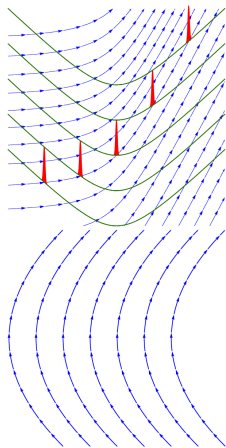
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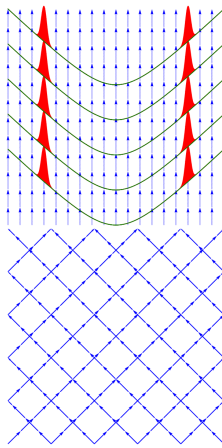
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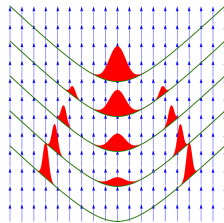
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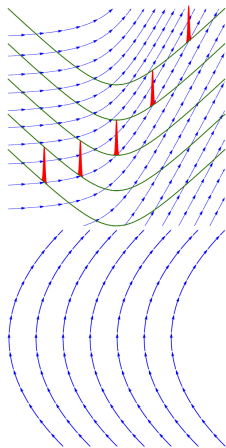
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Interacting fluid:
 $\mathcal{L}_r \phi \neq 0$.



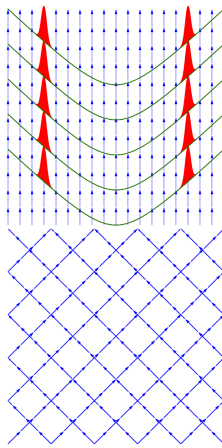
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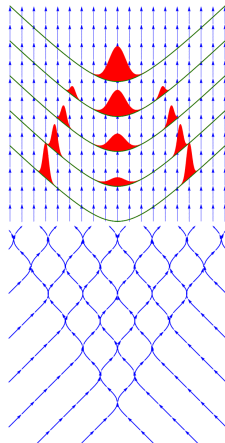
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Collisionless fluid:
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“Humppa”

Example: collisionless dust fluid

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- Apply Liouville equation:

$$0 = \nabla u^a = u^b \partial_b u^a + u^b N^a{}_b,$$

$$0 = \nabla_{\delta_a}(\rho u^a) = \partial_a(\rho u^a) + \frac{1}{2} \rho u^a g^{Fbc} \left(\partial_a g_{bc}^F - N^d{}_a \bar{\partial}_d g_{bc}^F \right).$$

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 - Mass density $\rho : M \rightarrow \mathbb{R}^+$.
 - Velocity $u : M \rightarrow O$.
- Particle density function:

$$\phi(x, y) \sim \rho(x) \delta_{S_x}(y, u(x)).$$

- Apply Liouville equation:

$$0 = \nabla u^a = u^b \partial_b u^a + u^b N^a{}_b,$$

$$0 = \nabla_{\delta_a}(\rho u^a) = \partial_a(\rho u^a) + \frac{1}{2} \rho u^a g^{Fbc} \left(\partial_a g_{bc}^F - N^d{}_a \bar{\partial}_d g_{bc}^F \right).$$

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⇒ Generalized (pressureless) Euler equations to Finsler geometry.

- Metric limit $F^2(x, y) = |g_{ab}(x)y^a y^b|$ yields Euler equations:

$$u^b \nabla_b u^a = 0, \quad \nabla_a(\rho u^a) = 0.$$

- Introduce suitable coordinates on TM :

$$t, r, \theta, \varphi, y^t, y^r, y^\theta, y^\varphi.$$

Cosmological symmetry

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- Most general Finsler function obeying cosmological symmetry:

$$F = F(t, y^t, w), \quad w^2 = \frac{(y^r)^2}{1 - kr^2} + r^2 \left((y^\theta)^2 + \sin^2 \theta (y^\varphi)^2 \right).$$

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- Homogeneity of Finsler function $F(t, y^t, w) = y^t \tilde{F}(t, w/y^t)$.
 - Introduce new coordinates: $\tilde{y} = y^t \tilde{F}(t, w/y^t)$, $\tilde{w} = w/y^t$.
- ⇒ Coordinates on observer space O with $\tilde{y} \equiv 1$.
- ⇒ Geometry function $\tilde{F}(t, \tilde{w})$ on O .

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Cosmological fluid dynamics

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- Example: collisionless dust fluid $\phi(x, y) \sim \rho(x) \delta_{S_x}(y, u(x))$:

$$u(t) = \frac{1}{\tilde{F}(t, 0)} \partial_t, \quad \partial_t \left(\rho(t) \sqrt{g^F(t, 0)} \right) = 0.$$

Outline

- 1 Motivation
- 2 Introduction to Finsler spacetimes
- 3 Finslerian description of fluids
- 4 Examples of fluids
- 5 Conclusion**

Conclusion

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 - Finsler spacetimes:
 - Define geometry by length functional.
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- Outlook:
 - Coupling of fluids to non-metric gravity theories.
 - Cosmological solutions with non-metric geometry.
 - Dark energy?
 - Inflation?
 - Extension of parameterized post-Newtonian formalism.

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