

Fluid dynamics on Finsler spacetimes

Manuel Hohmann

Laboratory of Theoretical Physics
Physics Institute
University of Tartu



DPG-Tagung Berlin – Session GR 15
19. März 2015

- **Fluids are everywhere:**
 - Perfect fluid (radiation, dust, dark matter. . .) - cosmology.
 - Maxwell-Boltzmann gas - atmospheres.
 - Plasma - stellar dynamics, primordial plasma.

- Fluids are everywhere:
 - Perfect fluid (radiation, dust, dark matter. . .) - cosmology.
 - Maxwell-Boltzmann gas - atmospheres.
 - Plasma - stellar dynamics, primordial plasma.
- Lift fluid dynamics to observer space:
 - All measurements are performed by observers.
 - Measurements depend on observer's frame (velocity).
 - Quantum gravity: possible non-tensorial velocity dependence.
 - Observer space: space of all physical velocities.
 - Fluids naturally modeled as densities on observer space.

- Fluids are everywhere:
 - Perfect fluid (radiation, dust, dark matter. . .) - cosmology.
 - Maxwell-Boltzmann gas - atmospheres.
 - Plasma - stellar dynamics, primordial plasma.
- Lift fluid dynamics to observer space:
 - All measurements are performed by observers.
 - Measurements depend on observer's frame (velocity).
 - Quantum gravity: possible non-tensorial velocity dependence.
 - Observer space: space of all physical velocities.
 - Fluids naturally modeled as densities on observer space.
- **Finsler spacetimes as observer space geometry:**
 - Finsler geometry of space widely used in physics.
 - Finsler geometry generalizes Riemannian geometry.
 - Finsler spacetimes are suitable backgrounds for physics.
 - Possible explanations of yet unexplained phenomena.

- Generalize metric length measure to Finsler function:

$$\tau = \int_{t_1}^{t_2} \sqrt{|g_{ab}(x(t))\dot{x}^a(t)\dot{x}^b(t)|} dt \rightsquigarrow \int_{t_1}^{t_2} F(x(t), \dot{x}(t)) dt .$$

- Generalize metric length measure to Finsler function:

$$\tau = \int_{t_1}^{t_2} \sqrt{|g_{ab}(x(t))\dot{x}^a(t)\dot{x}^b(t)|} dt \rightsquigarrow \int_{t_1}^{t_2} F(x(t), \dot{x}(t)) dt .$$

- Finsler function $F : TM \rightarrow \mathbb{R}^+$.
- Parametrization invariance requires homogeneity:

$$F(x, \lambda y) = \lambda F(x, y) \quad \forall \lambda > 0 .$$

- Generalize metric length measure to Finsler function:

$$\tau = \int_{t_1}^{t_2} \sqrt{|g_{ab}(x(t))\dot{x}^a(t)\dot{x}^b(t)|} dt \rightsquigarrow \int_{t_1}^{t_2} F(x(t), \dot{x}(t)) dt .$$

- Finsler function $F : TM \rightarrow \mathbb{R}^+$.
- Parametrization invariance requires homogeneity:

$$F(x, \lambda y) = \lambda F(x, y) \quad \forall \lambda > 0 .$$

- **Finsler spacetime** [C. Pfeifer, M. Wohlfarth '11]:
 - Length measure for tangent vectors.
 - Notion of timelike, lightlike, spacelike tangent vectors.
 - “Future unit timelike” vectors: physically allowed velocities.

- Generalize metric length measure to Finsler function:

$$\tau = \int_{t_1}^{t_2} \sqrt{|g_{ab}(x(t))\dot{x}^a(t)\dot{x}^b(t)|} dt \rightsquigarrow \int_{t_1}^{t_2} F(x(t), \dot{x}(t)) dt .$$

- Finsler function $F : TM \rightarrow \mathbb{R}^+$.
- Parametrization invariance requires homogeneity:

$$F(x, \lambda y) = \lambda F(x, y) \quad \forall \lambda > 0 .$$

- Finsler spacetime [C. Pfeifer, M. Wohlfarth '11]:
 - Length measure for tangent vectors.
 - Notion of timelike, lightlike, spacelike tangent vectors.
 - “Future unit timelike” vectors: physically allowed velocities.

⇒ Observer space $O \subset TM$ of allowed velocities.

Point mass dynamics on observer space

- Point mass follows curve $\gamma : \mathbb{R} \rightarrow M$ on spacetime M .
- γ is extremal curve of Finsler length measure:

$$\delta \int F(\gamma(t), \dot{\gamma}(t)) dt = 0.$$

Point mass dynamics on observer space

- Point mass follows curve $\gamma : \mathbb{R} \rightarrow M$ on spacetime M .
- γ is extremal curve of Finsler length measure:

$$\delta \int F(\gamma(t), \dot{\gamma}(t)) dt = 0.$$

- Canonical lift Γ of curve to tangent bundle TM :

$$\Gamma = (\gamma, \dot{\gamma}).$$

- Lift of geodesic equation to TM is first order differential equation:

$$\dot{\Gamma}(t) = \mathbf{S}(\Gamma(t)).$$

⇒ Solutions are integral curves of vector field \mathbf{S} on TM .

Point mass dynamics on observer space

- Point mass follows curve $\gamma : \mathbb{R} \rightarrow M$ on spacetime M .
- γ is extremal curve of Finsler length measure:

$$\delta \int F(\gamma(t), \dot{\gamma}(t)) dt = 0.$$

- Canonical lift Γ of curve to tangent bundle TM :

$$\Gamma = (\gamma, \dot{\gamma}).$$

- Lift of geodesic equation to TM is first order differential equation:

$$\dot{\Gamma}(t) = \mathbf{S}(\Gamma(t)).$$

⇒ Solutions are integral curves of vector field \mathbf{S} on TM .

- Physically allowed velocities: $\Gamma(t) \in O$.
- Restriction $\mathbf{r} = \mathbf{S}|_O$: Reeb vector field.

⇒ **Physical geodesics are integral curves of \mathbf{r} on O .**

From particles to fluids

- Model fluid by classical, relativistic particles:
 - Particles follow piecewise geodesic curves.
 - Endpoints of geodesics are interactions with other particles.

From particles to fluids

- Model fluid by classical, relativistic particles:
 - Particles follow piecewise geodesic curves.
 - Endpoints of geodesics are interactions with other particles.
- Particle measure: $\omega \in \Omega^6(\mathcal{O})$ unique 6-form such that:
 - ω non-degenerate on every hypersurface not tangent to \mathbf{r} .
 - $d\omega = 0$.

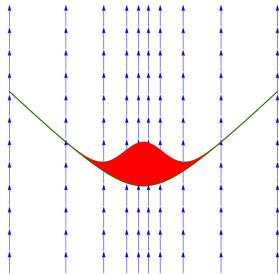
From particles to fluids

- Model fluid by classical, relativistic particles:
 - Particles follow piecewise geodesic curves.
 - Endpoints of geodesics are interactions with other particles.
- Particle measure: $\omega \in \Omega^6(\mathcal{O})$ unique 6-form such that:
 - ω non-degenerate on every hypersurface not tangent to \mathbf{r} .
 - $d\omega = 0$.
- Define one-particle distribution function $\phi : \mathcal{O} \rightarrow \mathbb{R}^+$ such that:

For every hypersurface $\sigma \subset \mathcal{O}$,

$$N[\sigma] = \int_{\sigma} \phi \omega$$

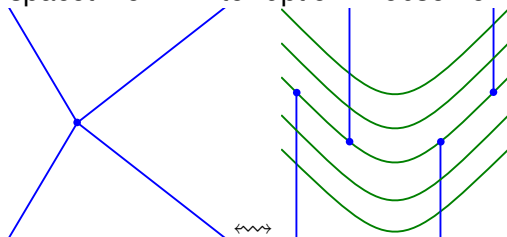
of **particle trajectories** through σ .



-
- Counting of particle trajectories respects hypersurface orientation.

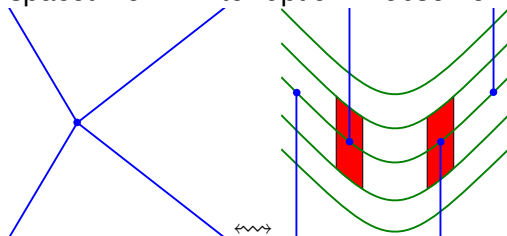
Collisions & the Liouville equation

- Collision in spacetime \leftrightarrow interruption in observer space.



Collisions & the Liouville equation

- Collision in spacetime \leftrightarrow interruption in observer space.



- For any open set $V \in \mathcal{O}$,

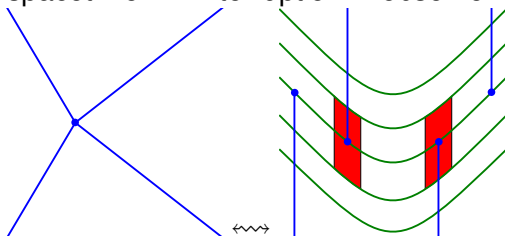
$$\int_{\partial V} \phi \omega = \int_V d(\phi \omega) = \int_V \mathcal{L}_r \phi \Sigma$$

of outbound trajectories - # of inbound trajectories.

\Rightarrow Collision density measured by $\mathcal{L}_r \phi$.

Collisions & the Liouville equation

- Collision in spacetime \leftrightarrow interruption in observer space.



- For any open set $V \in O$,

$$\int_{\partial V} \phi \omega = \int_V d(\phi \omega) = \int_V \mathcal{L}_{\mathbf{r}} \phi \Sigma$$

of outbound trajectories - # of inbound trajectories.

\Rightarrow Collision density measured by $\mathcal{L}_{\mathbf{r}} \phi$.

- **Collisionless fluid: trajectories have no endpoints, $\mathcal{L}_{\mathbf{r}} \phi = 0$.**

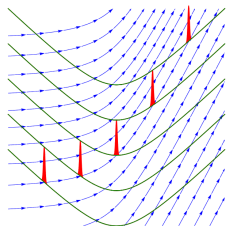
\Rightarrow Simple, first order equation of motion for collisionless fluid.

\Rightarrow ϕ is constant along integral curves of \mathbf{r} .

Examples of fluids

Geodesic dust fluid:

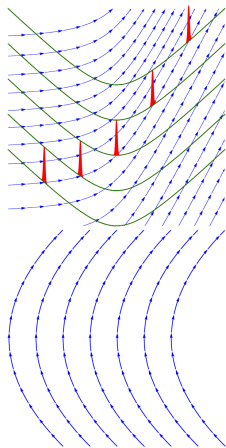
$$\phi(x, y) \sim \delta(y - u(x)).$$



Examples of fluids

Geodesic dust fluid:

$$\phi(x, y) \sim \delta(y - u(x)).$$

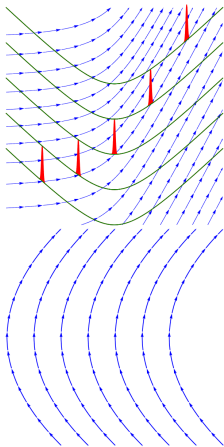


“Jenkka”

Examples of fluids

Geodesic dust fluid:

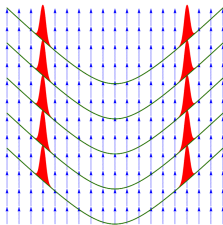
$$\phi(x, y) \sim \delta(y - u(x)).$$



“Jenkka”

Collisionless fluid:

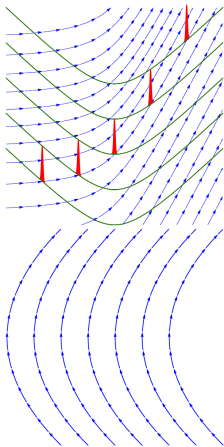
$$\mathcal{L}_r \phi = 0.$$



Examples of fluids

Geodesic dust fluid:

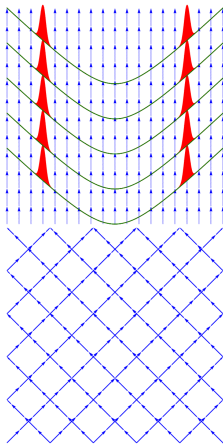
$$\phi(x, y) \sim \delta(y - u(x)).$$



“Jenkka”

Collisionless fluid:

$$\mathcal{L}_r \phi = 0.$$

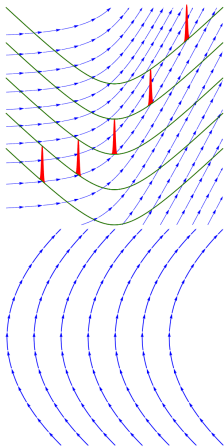


“Polkka”

Examples of fluids

Geodesic dust fluid:

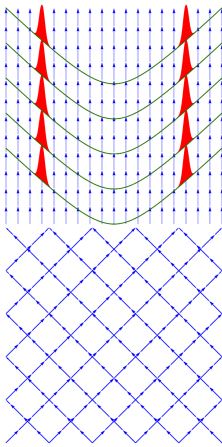
$$\phi(x, y) \sim \delta(y - u(x)).$$



“Jenkka”

Collisionless fluid:

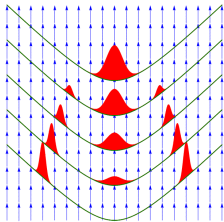
$$\mathcal{L}_r \phi = 0.$$



“Polkka”

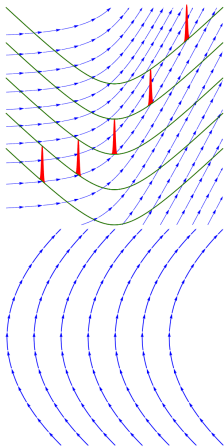
Interacting fluid:

$$\mathcal{L}_r \phi \neq 0.$$



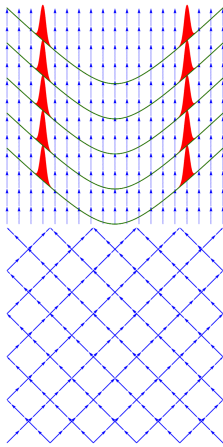
Examples of fluids

Geodesic dust fluid:
 $\phi(x, y) \sim \delta(y - u(x))$.



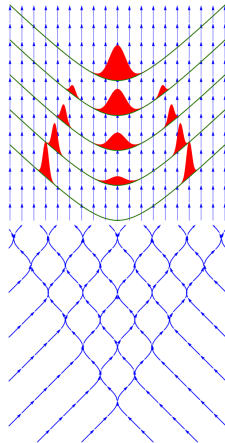
“Jenkka”

Collisionless fluid:
 $\mathcal{L}_r \phi = 0$.



“Polkka”

Interacting fluid:
 $\mathcal{L}_r \phi \neq 0$.



“Humppa”

Example: dust fluid in Finsler geometry

- Classical dust with density $\rho(x)$ and velocity $u^a(x)$.
- One-particle distribution function on O :

$$\phi(\hat{x}, \theta) = \frac{1}{m} \rho(\hat{x}) \frac{\delta(\theta - v(\hat{x}))}{\sqrt{\det h^F(\hat{x}, \theta)}}.$$

Example: dust fluid in Finsler geometry

- Classical dust with density $\rho(x)$ and velocity $u^a(x)$.
- One-particle distribution function on O :

$$\phi(\hat{x}, \theta) = \frac{1}{m} \rho(\hat{x}) \frac{\delta(\theta - v(\hat{x}))}{\sqrt{\det h^F(\hat{x}, \theta)}}.$$

- **No collisions \Rightarrow Liouville equation $\mathcal{L}_r \phi = 0$.**

\Rightarrow Equations of motion for ρ and u^a :

$$\nabla u^a = 0 \quad \text{and} \quad \nabla_{\delta_a}(\rho u^a) = 0.$$

- Dynamical covariant derivative ∇ .
- Horizontal part of Cartan linear connection ∇_{δ_a} .

Example: dust fluid in Finsler geometry

- Classical dust with density $\rho(x)$ and velocity $u^a(x)$.
- One-particle distribution function on O :

$$\phi(\hat{x}, \theta) = \frac{1}{m} \rho(\hat{x}) \frac{\delta(\theta - v(\hat{x}))}{\sqrt{\det h^F(\hat{x}, \theta)}}.$$

- No collisions \Rightarrow Liouville equation $\mathcal{L}_r \phi = 0$.

\Rightarrow Equations of motion for ρ and u^a :

$$\nabla u^a = 0 \quad \text{and} \quad \nabla_{\delta_a}(\rho u^a) = 0.$$

- Dynamical covariant derivative ∇ .
- Horizontal part of Cartan linear connection ∇_{δ_a} .
- Metric background geometry $F(x, y) = \sqrt{|g_{ab}(x)y^a y^b|}$:

$$u^b \nabla_b u^a = 0 \quad \text{and} \quad \nabla_a(\rho u^a) = 0.$$

\Rightarrow Well-known Euler equations of fluid dynamics.

Fluids with cosmological symmetry

- Most general cosmological Finsler function $F(\hat{t}, \hat{y}, \hat{w})$.
 - Cosmological time \hat{t} .
 - Velocity component \hat{y} in \hat{t} -direction.
 - Velocity component \hat{w} perpendicular to \hat{t} -direction.

Fluids with cosmological symmetry

- Most general cosmological Finsler function $F(\hat{t}, \hat{y}, \hat{w})$.
 - Cosmological time \hat{t} .
 - Velocity component \hat{y} in \hat{t} -direction.
 - Velocity component \hat{w} perpendicular to \hat{t} -direction.
- Homogeneity: F determined by \tilde{F} as

$$F(\hat{t}, \hat{y}, \hat{w}) = \hat{y} \tilde{F}(\hat{t}, \hat{w}/\hat{y}).$$

- Observer space O with $\hat{y} \tilde{F}(\hat{t}, \hat{w}/\hat{y}) = 1$.

Fluids with cosmological symmetry

- Most general cosmological Finsler function $F(\hat{t}, \hat{y}, \hat{w})$.
 - Cosmological time \hat{t} .
 - Velocity component \hat{y} in \hat{t} -direction.
 - Velocity component \hat{w} perpendicular to \hat{t} -direction.
- Homogeneity: F determined by \tilde{F} as

$$F(\hat{t}, \hat{y}, \hat{w}) = \hat{y} \tilde{F}(\hat{t}, \hat{w}/\hat{y}).$$

- Observer space O with $\hat{y} \tilde{F}(\hat{t}, \hat{w}/\hat{y}) = 1$.
- Most general fluid $\phi(\hat{t}, \hat{w}/\hat{y})$ with cosmological symmetry.
- **Liouville equation** $\mathcal{L}_r \phi = 0$:

$$\tilde{F}_{ww} \phi_t = \tilde{F}_{tw} \phi_w.$$

Fluids with cosmological symmetry

- Most general cosmological Finsler function $F(\hat{t}, \hat{y}, \hat{w})$.
 - Cosmological time \hat{t} .
 - Velocity component \hat{y} in \hat{t} -direction.
 - Velocity component \hat{w} perpendicular to \hat{t} -direction.
- Homogeneity: F determined by \tilde{F} as

$$F(\hat{t}, \hat{y}, \hat{w}) = \hat{y} \tilde{F}(\hat{t}, \hat{w}/\hat{y}).$$

- Observer space O with $\hat{y} \tilde{F}(\hat{t}, \hat{w}/\hat{y}) = 1$.
- Most general fluid $\phi(\hat{t}, \hat{w}/\hat{y})$ with cosmological symmetry.
- Liouville equation $\mathcal{L}_r \phi = 0$:

$$\tilde{F}_{ww} \phi_t = \tilde{F}_{tw} \phi_w.$$

- Robertson-Walker metric: $\tilde{F} = \sqrt{1 - a^2(\hat{t}) \hat{w}^2 / \hat{y}^2}$:

$$\phi_t = -\frac{\hat{w}}{\hat{y}} \left(\frac{\hat{w}^2}{\hat{y}^2} a^2 - 2 \right) \frac{\dot{a}}{a} \phi_w.$$

- Basic idea:
 - Model fluids by particle trajectories.
 - Lift trajectories from spacetime to observer space.
 - Describe geometry of observer space using Finsler geometry.
 - Measure particle density by distribution function.
 - Derive fluid dynamics from geodesic motion.

- Basic idea:
 - Model fluids by particle trajectories.
 - Lift trajectories from spacetime to observer space.
 - Describe geometry of observer space using Finsler geometry.
 - Measure particle density by distribution function.
 - Derive fluid dynamics from geodesic motion.
- Presented examples:
 - Classical dust fluid on Finsler spacetime.
 - Most general cosmological fluid on Finsler spacetime.

- Basic idea:
 - Model fluids by particle trajectories.
 - Lift trajectories from spacetime to observer space.
 - Describe geometry of observer space using Finsler geometry.
 - Measure particle density by distribution function.
 - Derive fluid dynamics from geodesic motion.
- Presented examples:
 - Classical dust fluid on Finsler spacetime.
 - Most general cosmological fluid on Finsler spacetime.
- Future research goals:
 - Coupling of fluids to non-metric gravity theories.
 - Cosmological solutions of gravity with non-metric geometry.
 - Extension of parameterized post-Newtonian formalism.

- Kinetic theory on the tangent bundle:
 - J. Ehlers, in: “General Relativity and Cosmology”, pp 1–70, Academic Press, New York / London, 1971.
 - O. Sarbach and T. Zannias, AIP Conf. Proc. **1548** (2013) 134 [arXiv:1303.2899 [gr-qc]].
 - O. Sarbach and T. Zannias, Class. Quant. Grav. **31** (2014) 085013 [arXiv:1309.2036 [gr-qc]].
- Finsler spacetimes:
 - C. Pfeifer and M. N. R. Wohlfarth, Phys. Rev. D **84** (2011) 044039 [arXiv:1104.1079 [gr-qc]].
 - C. Pfeifer and M. N. R. Wohlfarth, Phys. Rev. D **85** (2012) 064009 [arXiv:1112.5641 [gr-qc]].
 - MH, in: “Mathematical structures of the Universe”, pp 13–55, Copernicus Center Press, Krakow, 2014.